



BEAM DYNAMICS OF IONIZATION COOLING CHANNELS

Kwang-Je Kim and Chun-Xi Wang
The University of Chicago and ANL

MUTAC Meeting
Brookhaven National Laboratory
June 15-16, 2000

TOPICS

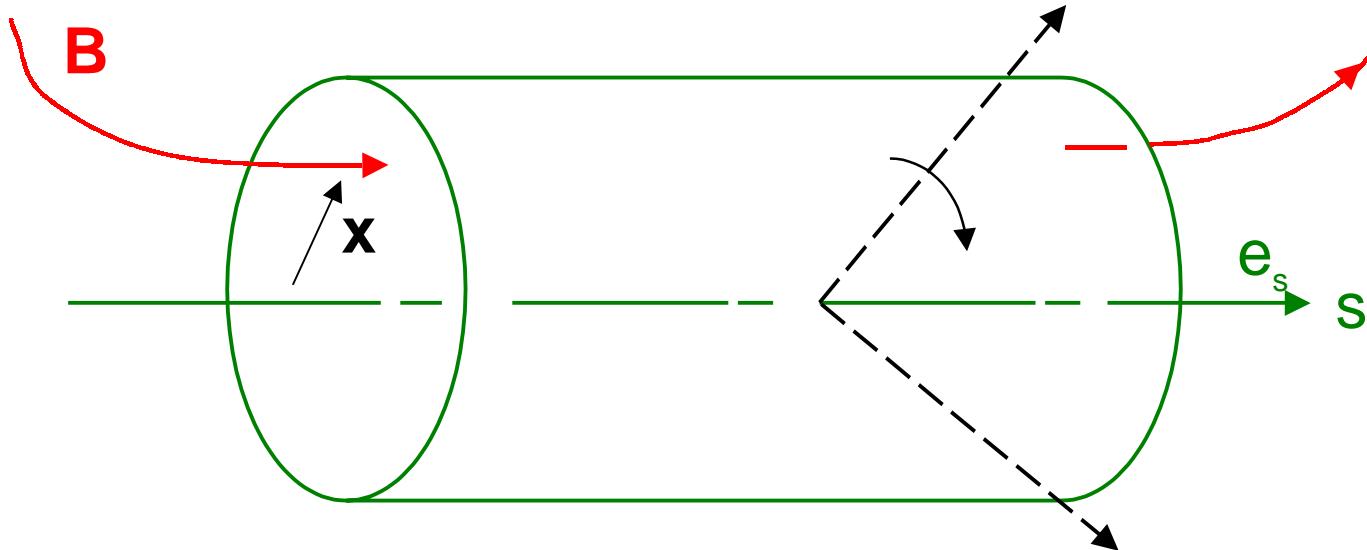
- Linear Transverse Cooling Dynamics

Penn & Wurtele (PRL), Kim & Wang (PRL), Kim & Wang (MUCool)

- Orbit Stability and Energy Acceptance

Wang & Kim (MUCool)

SOLENOIDAL FOCUSING



$B(s)$: solenoidal field on s -axis

Simplest 3-D field satisfying $\nabla \cdot \mathbf{B} = 0, \nabla \times \mathbf{B} = 0$

$$\mathbf{B}(s, \mathbf{x}_\perp) = B(s) \mathbf{e}_s - \frac{1}{2} \left(\frac{dB}{ds} \right) \mathbf{x}_\perp + \text{non-linear terms}$$

EQUATION OF MOTION

- Lab frame:

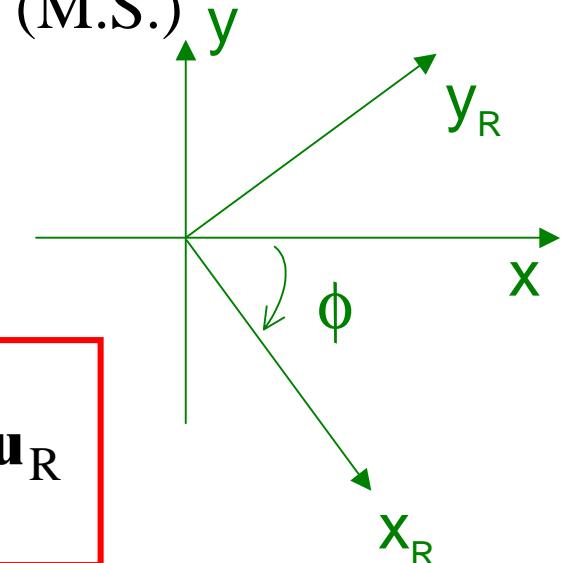
$$\frac{d}{ds} P_s \frac{dx_{\perp}}{ds} = -qB\mathbf{e}_s \times \frac{dx}{ds} - \frac{q}{2} B' \mathbf{e}_s \times \mathbf{x} - \eta P_s \frac{dx}{ds} + P_s \mathbf{u}$$

$$\eta = \left. \frac{1}{P_s} \frac{dP_s}{ds} \right|_{\text{loss}} ; \quad \mathbf{u} : \text{stochastic excitation (M.S.)}$$

- Larmor frame:

$$\frac{d}{ds} P_s \frac{dx_R}{ds} = -P_s \kappa^2 \mathbf{x}_R - P_s \eta \left(\frac{dx_R}{ds} - \kappa \mathbf{e}_s \times \mathbf{x}_R \right) - P_s \mathbf{u}_R$$

$$\frac{d\phi}{ds} = \kappa = \frac{qB}{2P_s}$$



ENVELOPE EQUATIONS IN LARMOR FRAME

Phase space variables: $\{\mathbf{x}, \mathbf{x}'\} = \{x, y, x', y'\}$

$$\frac{d\mathbf{x}}{ds} = \mathbf{x}'$$

$$\frac{d\mathbf{x}'}{ds} = -\kappa^2 \mathbf{x} - \eta(\mathbf{x}' - \kappa \mathbf{e}_s \times \mathbf{x}) - \mathbf{u}$$

$$\langle \mathbf{x}^2 \rangle = \varepsilon \beta \quad \varepsilon = \sqrt{\langle \mathbf{x}^2 \rangle \langle \mathbf{x}'^2 \rangle - \langle \mathbf{x} \cdot \mathbf{x}' \rangle}$$

$$\langle \mathbf{x}'^2 \rangle = \varepsilon \gamma \quad \beta \gamma - \alpha^2 = 1$$

$$\langle \mathbf{x} \cdot \mathbf{x}' \rangle = -\varepsilon \alpha$$

$$\mathbf{e}_s \cdot \langle \mathbf{x} \times \mathbf{x}' \rangle = \langle xy' - yx' \rangle = \mathbf{L}$$

Excitation $\chi = n \langle \mathbf{u}^2 \rangle = \left(\frac{13.6 \text{ MeV}}{pv} \right)^2 \frac{1}{L_{\text{Rad}}}$

EVOLUTION OF "LATTICE FUNCTIONS"

$$\frac{d\beta}{ds} + 2\alpha = \left(\eta\beta - \frac{\eta\kappa L + \chi}{\varepsilon}\beta^2 \right)$$

$$\frac{d\alpha}{ds} + \frac{1+\alpha^2}{\beta} - \eta^2\beta = \left(-\frac{\eta\kappa L + \chi}{\varepsilon}\beta\alpha \right)$$

R.H.S. are small!

$$\Rightarrow \frac{1}{2}\beta'' + \kappa^2\beta - \frac{1}{\beta} \left(1 + \frac{\beta'^2}{4} \right) = 0$$

EVOLUTION OF "BEAM" PARAMETERS:

$$\frac{d}{ds} \begin{pmatrix} \varepsilon \\ L \\ \Psi \end{pmatrix} = - \begin{pmatrix} \eta & -\eta\kappa\beta \\ -\eta\kappa\beta & \eta \end{pmatrix} \begin{pmatrix} \varepsilon \\ L \end{pmatrix} + \begin{pmatrix} \beta\chi \\ 0 \end{pmatrix} \quad \Xi$$

General Solution:

$$\Psi(s) = e^{-\hat{\Gamma}(s)} \Psi(0) + e^{-\hat{\Gamma}(s)} \int_0^s d\bar{s} e^{-\hat{\Gamma}(\bar{s})} \Xi(\bar{s})$$

$$\hat{\Gamma}(s) \equiv \int_0^s d\bar{s} \hat{M}(\bar{s}) = \begin{pmatrix} \varsigma_1 & -\varsigma_2(s) \\ -\varsigma_2(s) & \varsigma_1(s) \end{pmatrix}$$

$$\varsigma_1(s) \equiv \int_0^s d\bar{s} \eta(\bar{s}), \quad \varsigma_2(s) \equiv \int_0^s d\bar{s} \eta(\bar{s}) \kappa(\bar{s}) \beta(\bar{s})$$

PERIODIC COOLING CHANNELS

$$\psi(m\lambda) = e^{-m\Gamma(\lambda)} \psi(0)$$

$$+ \left\{ e^{-\Gamma(\lambda)} + e^{-2\Gamma(\lambda)} + \dots + e^{-(m-1)\Gamma(\lambda)} \right\} e^{-\Gamma(\lambda)} \int_0^\lambda d\bar{s} e^{\Gamma(s)} \begin{pmatrix} \beta\chi \\ 0 \end{pmatrix}$$

$$\psi(0) = \begin{pmatrix} \varepsilon_0 \\ L_0 \end{pmatrix} \quad \psi(m\lambda)$$

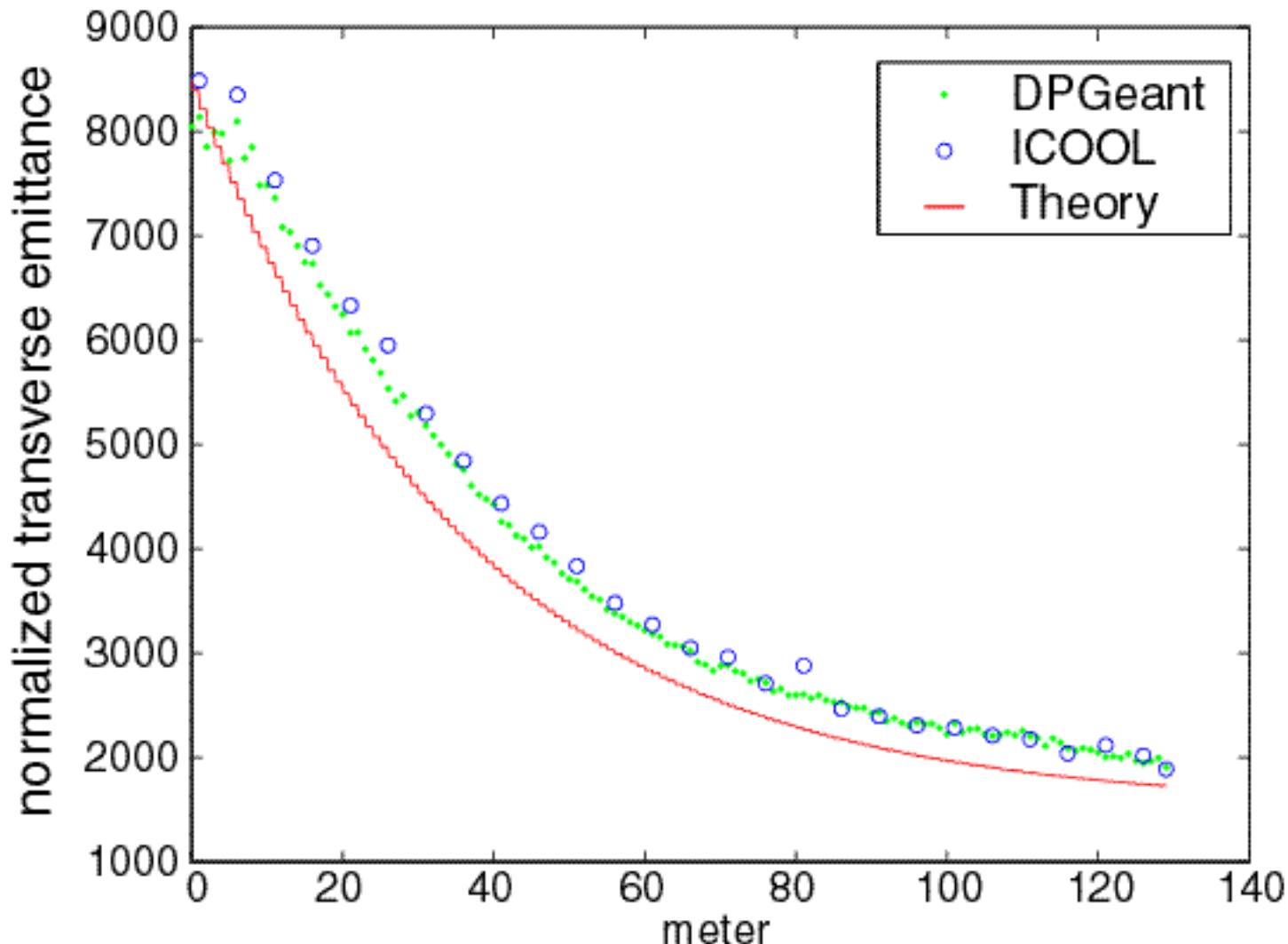


COMPARISON OF TRANSVERSE EMITTANCE (COOLING FACTOR) CALCULATIONS FOR MUON COOLING CHANNELS

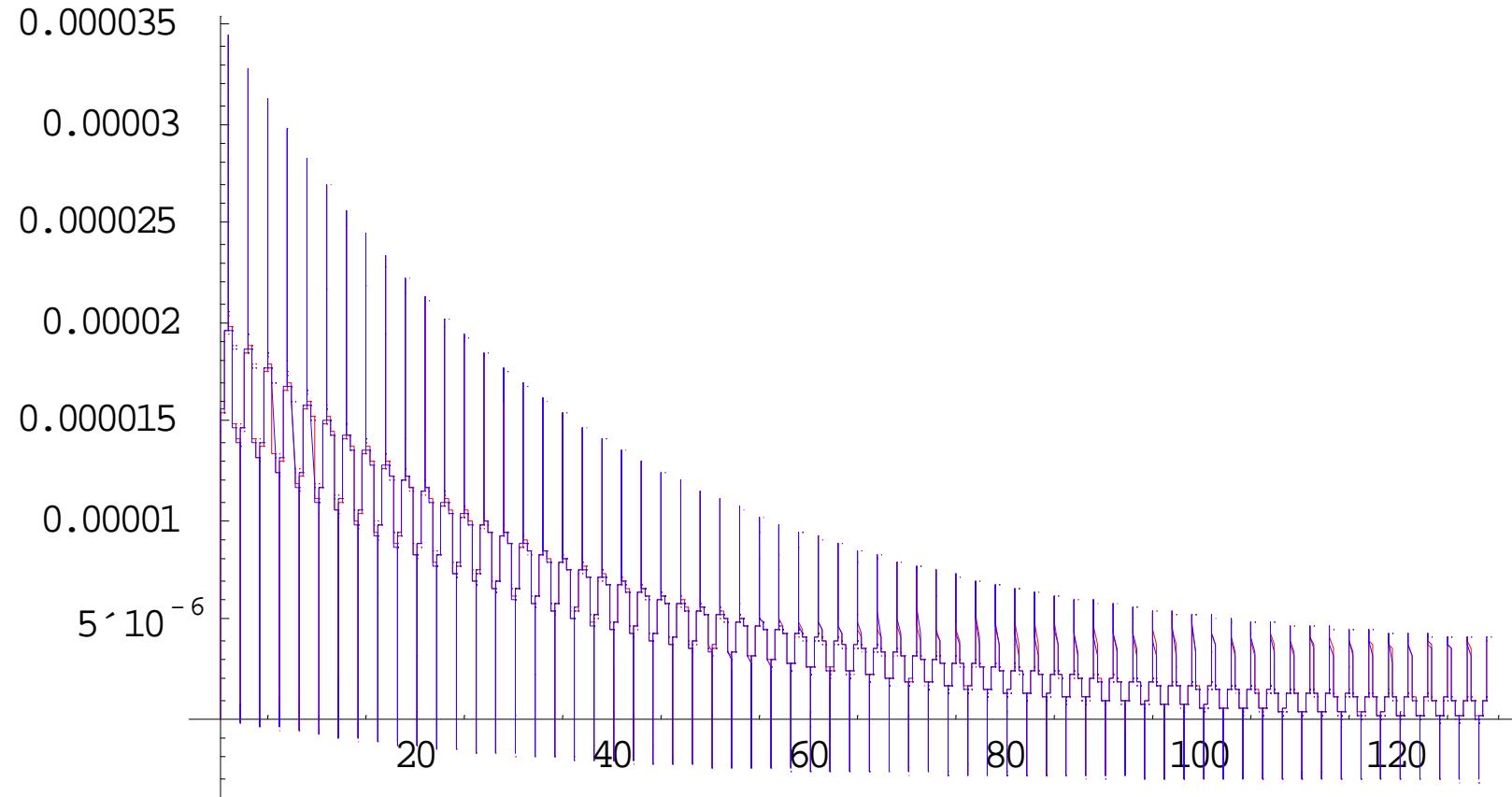
	3.4T FOFO (1)	3.4T FOFO (2)	DFOFO
ε^∞	0.19	0.27	0.080
$\varepsilon(m\lambda)$	0.21	0.34	0.095
ICOOL	0.22	0.34	0.10
DPGeant	0.24	0.37	N/A

$$\varepsilon^\infty \cong \frac{\int_0^\lambda d\bar{s} \beta(\bar{s}) \chi(\bar{s})}{\int_0^\lambda d\bar{s} \eta(\bar{s})}$$

COMPARISON OF COOLING PERFORMANCE CALCULATION



CANONICAL ANGULAR MOMENTUM IN A FOFO CHANNEL



ANGULAR MOMENTUM MODIFIES NEUFFER EQUATION!

$$\frac{d\epsilon}{ds} = -\eta\epsilon + \beta\chi + \eta\kappa\beta L$$

$$\frac{dL}{ds} = -\eta L + \eta\kappa\beta\epsilon$$

Effect of angular momentum

Neuffer

Equilibrium Emittance:

$$\epsilon^\infty = \frac{\int_0^\lambda ds \beta(s) \chi(s)}{\int_0^\lambda ds \eta(s)} \mathcal{F}$$

$$\mathcal{F} = 1 \Bigg/ \left(1 - \left[\int_0^\lambda ds \eta \kappa \beta \right] \Big/ \left[\int_0^\lambda ds \eta \right]^2 \right)$$

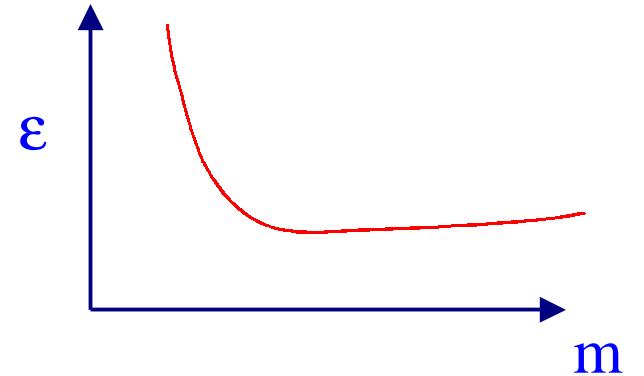
$\mathcal{F} = 1$ for alternating solenoids [R. Palmer]

MINIMUM EMITTANCE FOR SINGLE FLIP

Constant magnetic field

$$\kappa\beta = 1, \quad \mathcal{F} \rightarrow \infty$$

$$\varepsilon^{(m)} \rightarrow m \text{ (linearly divergent)}$$



Minimum emittance for $\varepsilon^\infty \ll \varepsilon(0)$

$$\varepsilon_{\min} \approx \varepsilon^\infty \left[\frac{3}{4} + \frac{1}{2} \ln \left(\frac{4\varepsilon(0)}{\varepsilon^\infty} \right) \right]$$

(ε^∞ = equilibrium emittance for alternating case with the same β)

Single Flip Cooling Channel: Principle

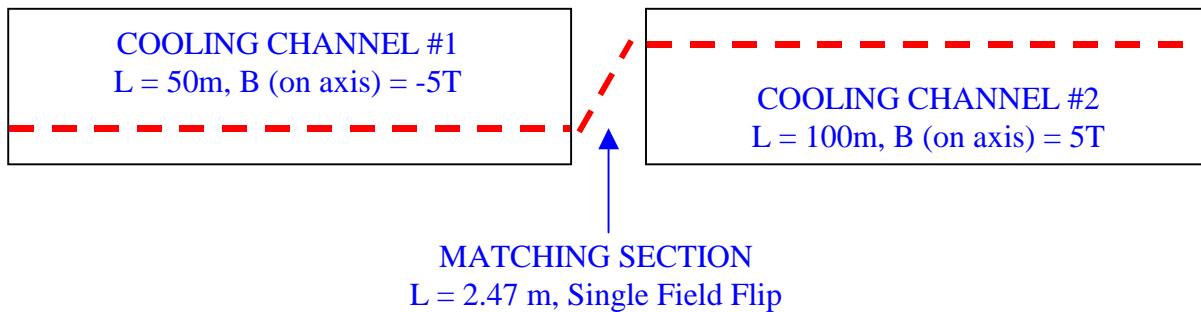


Figure 1: Schematic Layout of the Single Flip Channel

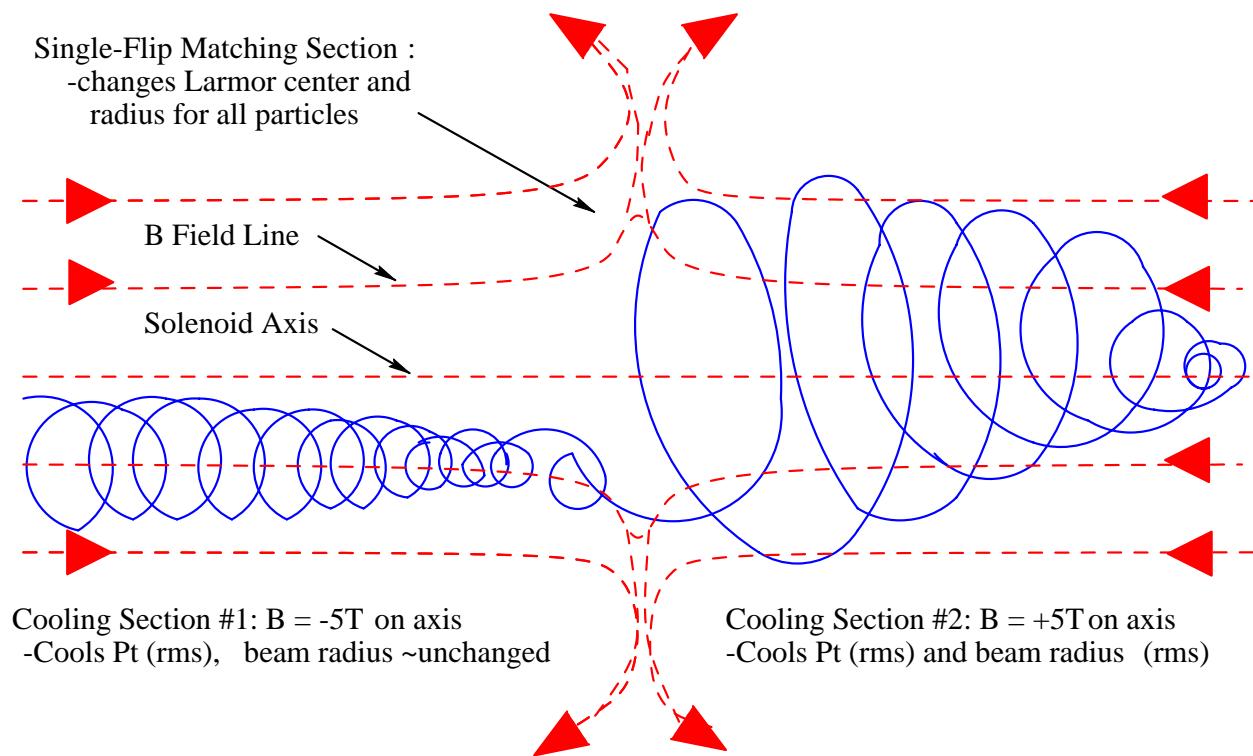
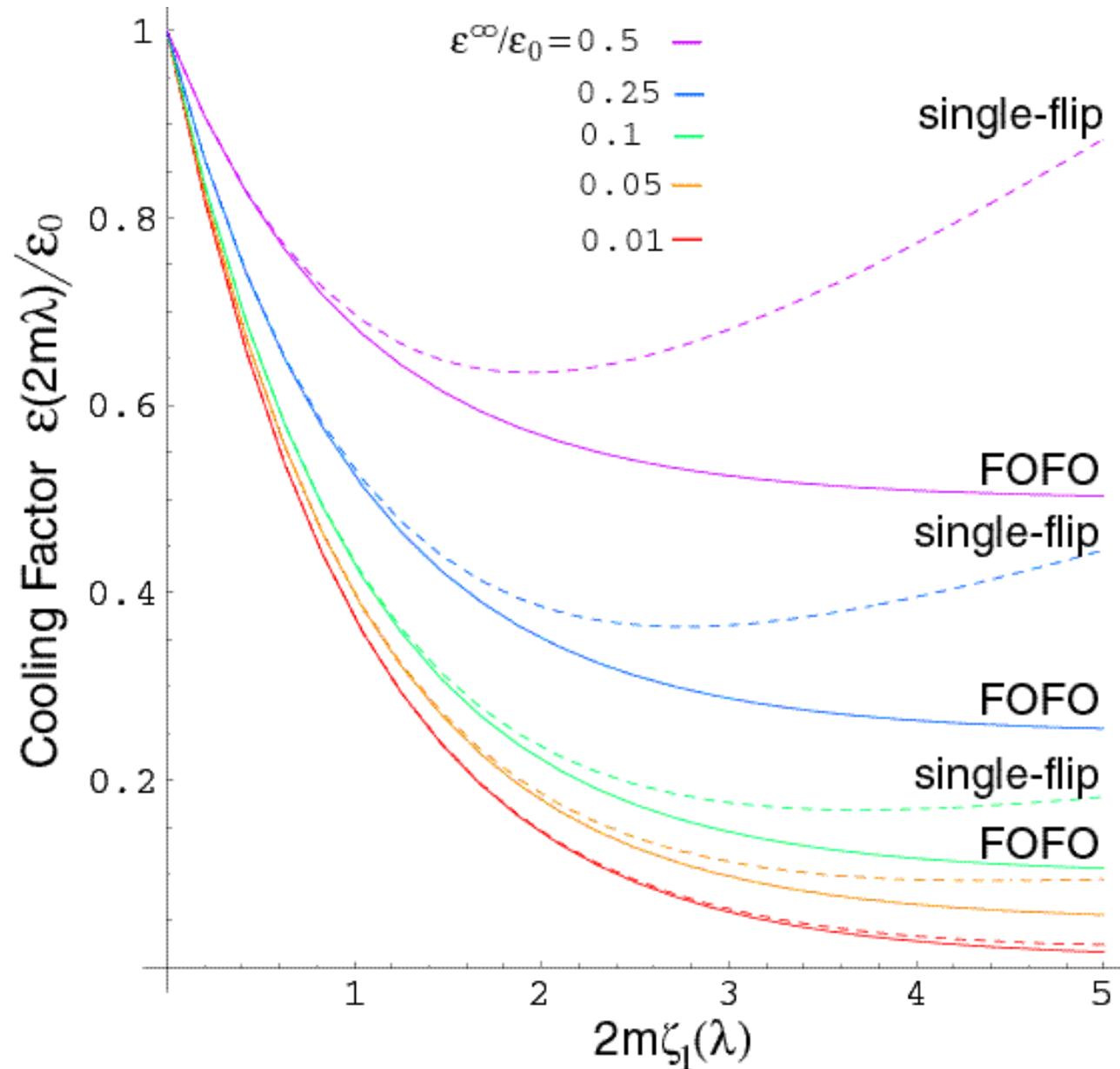


Figure 2: Particle Motion Cartoon in the Single Flip Channel

COMPARISON OF FOFO AND SINGLE-FLIP COOLING CHANNEL



HILL'S EQUATION

$$\frac{d^2 \mathbf{x}}{ds^2} + \kappa^2(s) \mathbf{x}(s) = 0$$

$$\kappa^2(s) = \theta_0 + 2\theta_1 \cos 2s + 2\theta_2 \cos 4s + \dots; \quad \theta_0 = \left(\frac{qL}{2\pi P_s} \right)^2 \langle B_s^2 \rangle$$

G.W. Hill "On the Part of the Motion of the Lunar Perigee
which is a Function of Mean Motion of Sun &
Moon" (1886)

Whittaker & Watson
Magnus & Winkler

STABILITY CRITERIA

$$\text{Tr } R(\pi) = \Delta = 2 - 4 \sin^2 \left(\frac{\pi}{2} \sqrt{\theta_0} \right) \cdot D$$

$$D = \left\| \delta_{nm} + \frac{\theta_{n-m}}{\theta_0 - 4n^2} \right\| \quad (\text{Hill's determinant})$$

$$\Delta = 2 \cos(\sqrt{\theta_0} \pi) + \frac{\pi \sin \sqrt{\theta_0} \pi}{2\sqrt{\theta_0}} \sum_{n=1}^{\infty} \frac{|\theta_n|^2}{\theta_0 - n^2} + \dots$$

Stability $|\Delta| \leq 2$.

Stopband centered at $\sqrt{\theta_0} \approx n \left(1 + \frac{5}{16} \frac{\theta_n^2}{\theta_0^2} \right)$

Width: $\Delta \sqrt{\theta_0} \approx n \frac{\theta_n}{\theta_0}$
$$\left[\frac{\Delta P_z}{P_z \approx \frac{\theta_n}{\theta_0}} \right]$$

FOFO

$$B(s) = B_0 \sin s$$

$$B^2(s) = B_0^2 (1 - \cos 2s)$$

$\theta_0 = \theta_1 \rightarrow$ "maximally unstable"

To reduce θ_1

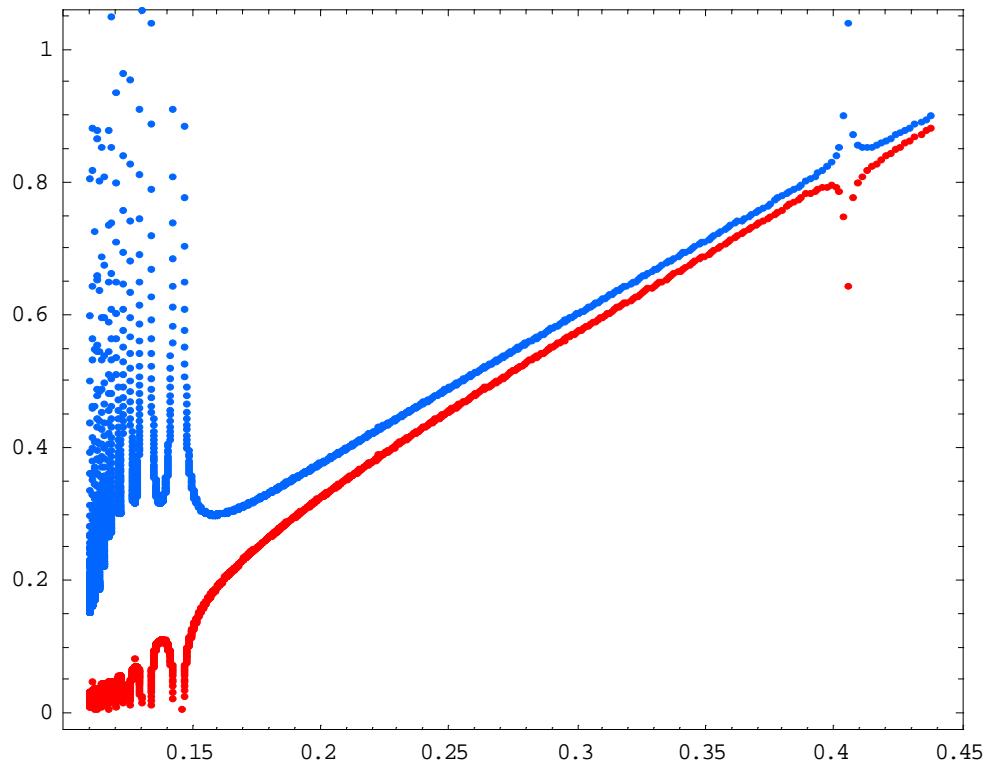
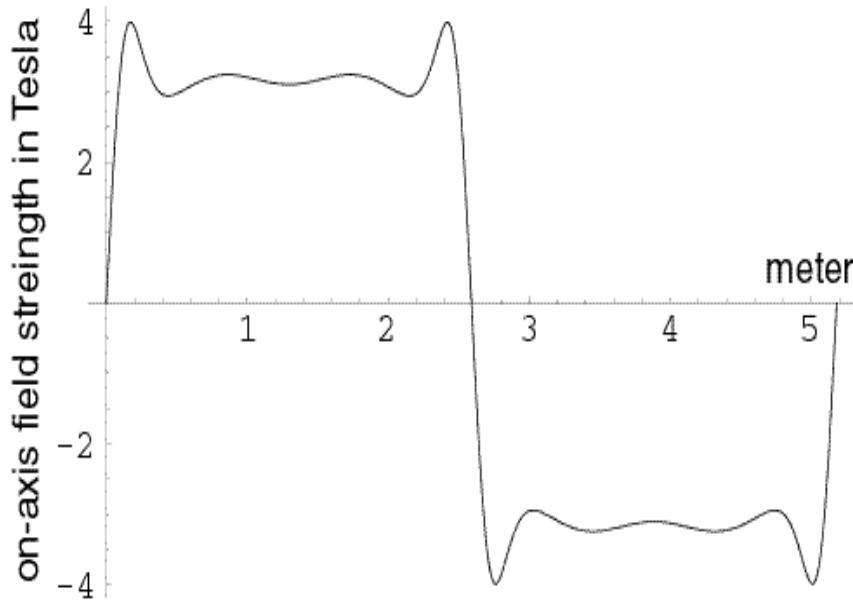
$$B(s) = B_0 (\sin s + a \sin 3s + \dots)$$

$$B^2(s) = B_0^2 (\sin^2 s + a^2 \sin^2 3s + 2a \sin s \sin 3s + \dots)$$

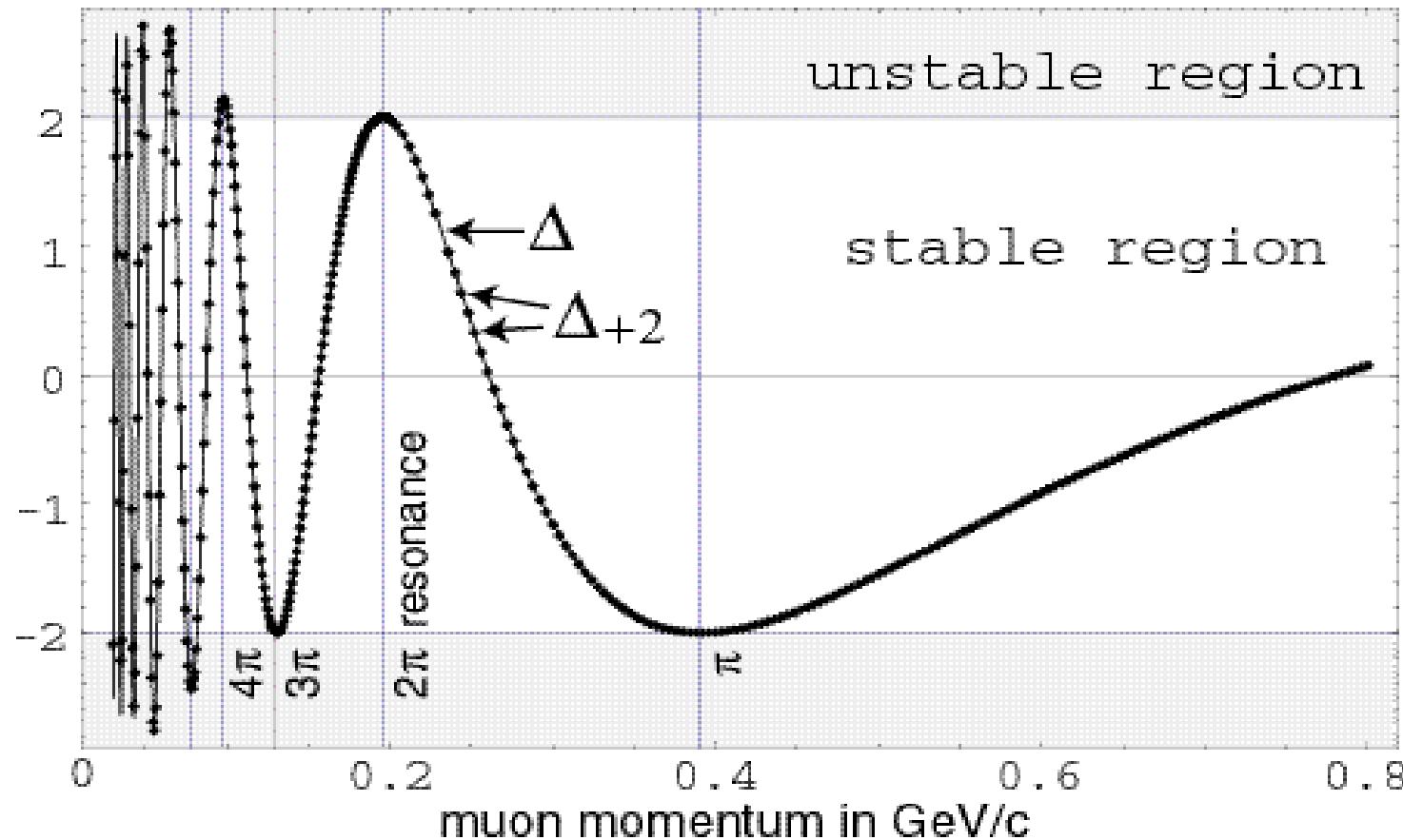
$$B_0^2 \left(\frac{1 - \cos 2s}{2} + \frac{a^2 (1 - \cos 6s)}{2} + a(\cos 2s - \cos 4s) + \dots \right)$$

Choose $a = \frac{1}{2}$

V. BALBEKOV'S FAST-FLIP COOLING CHANNEL

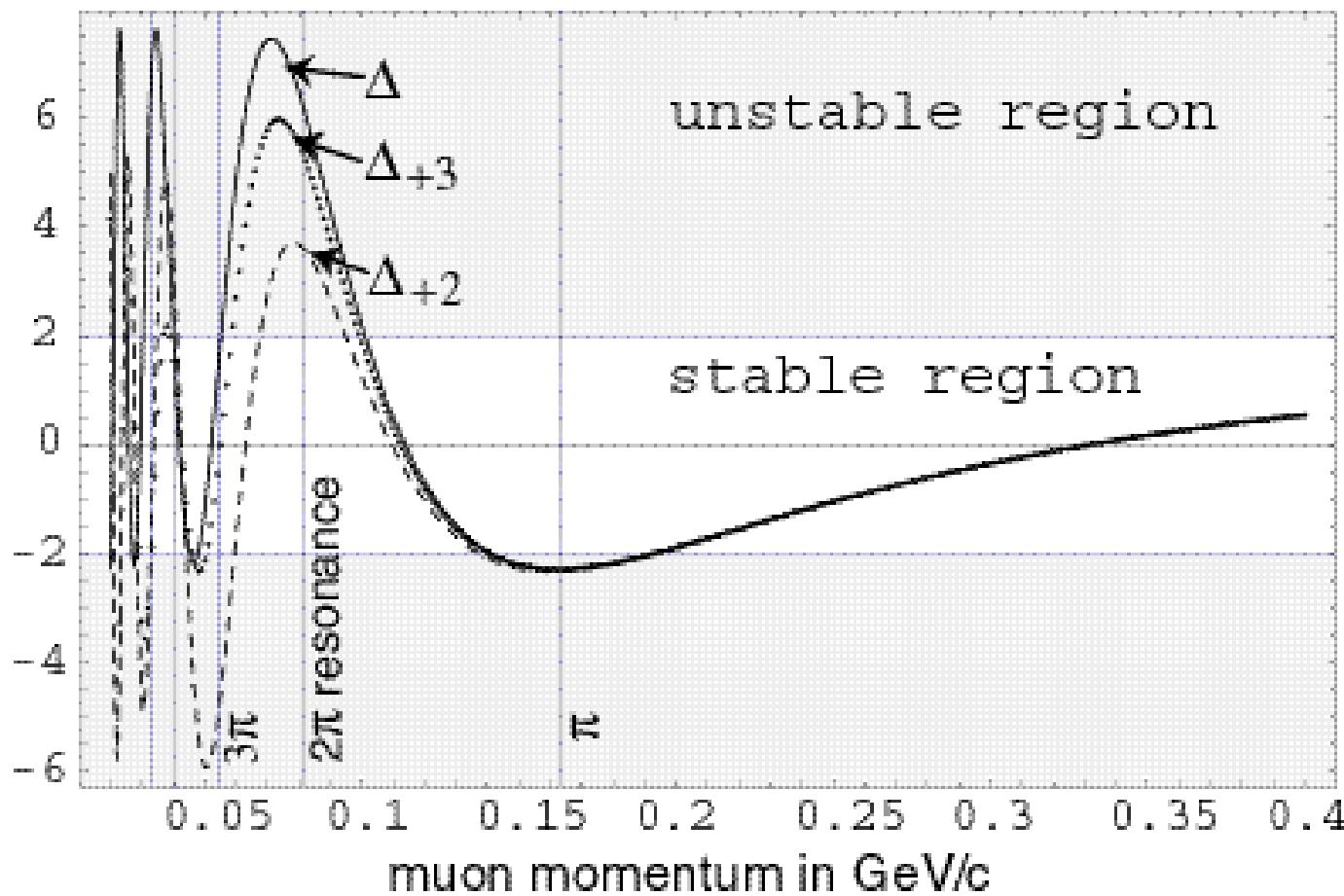


RESONANCE SUPPRESSION IN V. BALBEKOV'S CASE



At 200 MeV, $\{\kappa_0, \kappa_1, \dots, \kappa_{14}\} =$
 $\{3.8186, -0.0039, -0.0047, -0.0060, -0.2239, -0.3203, -0.3345, -0.2976,$
 $-0.2383, -0.1759, -0.1212, -0.0783, -0.0477, -0.0272, -0.0146\}$

APPLICATION TO G. PENN'S SUPERFOFO EXAMPLE



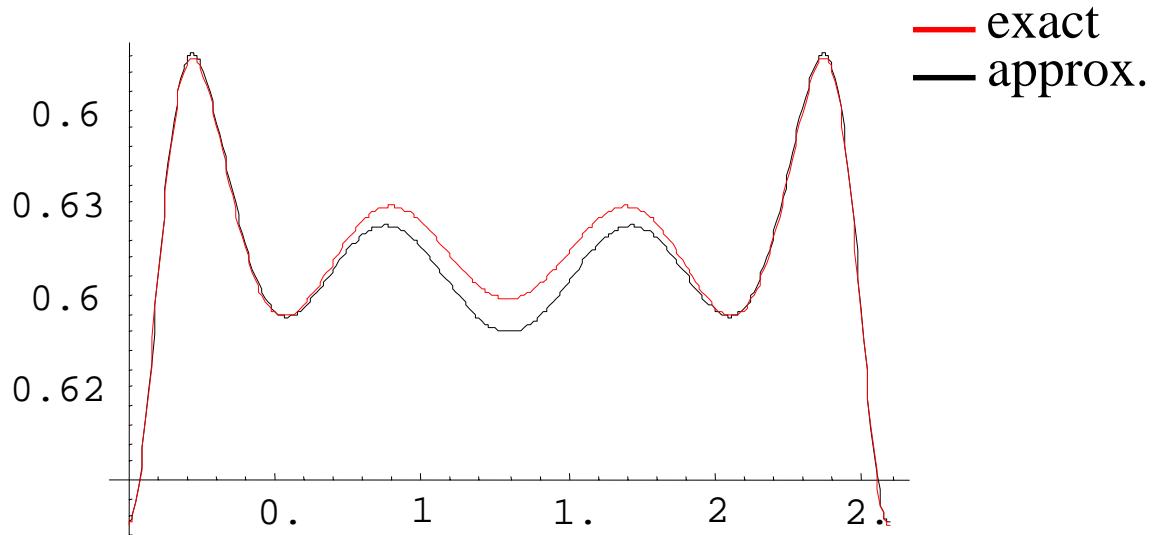
LEARNING TO OPTIMIZE (MINIMIZE) BETA FUNCTION

$$\beta(s) = \frac{L \sin(\sqrt{\theta_0} \pi)}{\pi \sqrt{\theta_0} (\sin \psi)} \left\{ 1 + \sum_{n=1}^{\infty} \frac{\operatorname{Re}[\theta_n e^{i2ns}]}{n^2 - \theta_0} + \dots \right\}$$

- Minimize L
- Large phase advance ψ
- Maximize $\theta_0 = \frac{Lq}{2\pi P_s} \sqrt{\langle B_s^2 \rangle}$ (higher pass band)
- Large θ_n to cancel 1

*All these within the stability limit!

V. Balbekov fast field flip channel at P=300MeV



SuperFOFO example of G. Penn at P=200 MeV

