

*Naturalness Principles
Supersymmetry and
Tenable Physics Motivations for
HEMC's*

Greg Anderson

Northwestern University

September 26, 1999

HEMC'99, Sept, 1999

Outline:

1. Naturalness and Superpartner Masses

- Why are superpartners "heavy"
- Fundamental Scales
- The naturalness problem
- 3 Solutions
- Measures of Fine-tuning
- Fine tuning in the MSSM
- Expected values of superpartner masses
- Variations on the MSSM

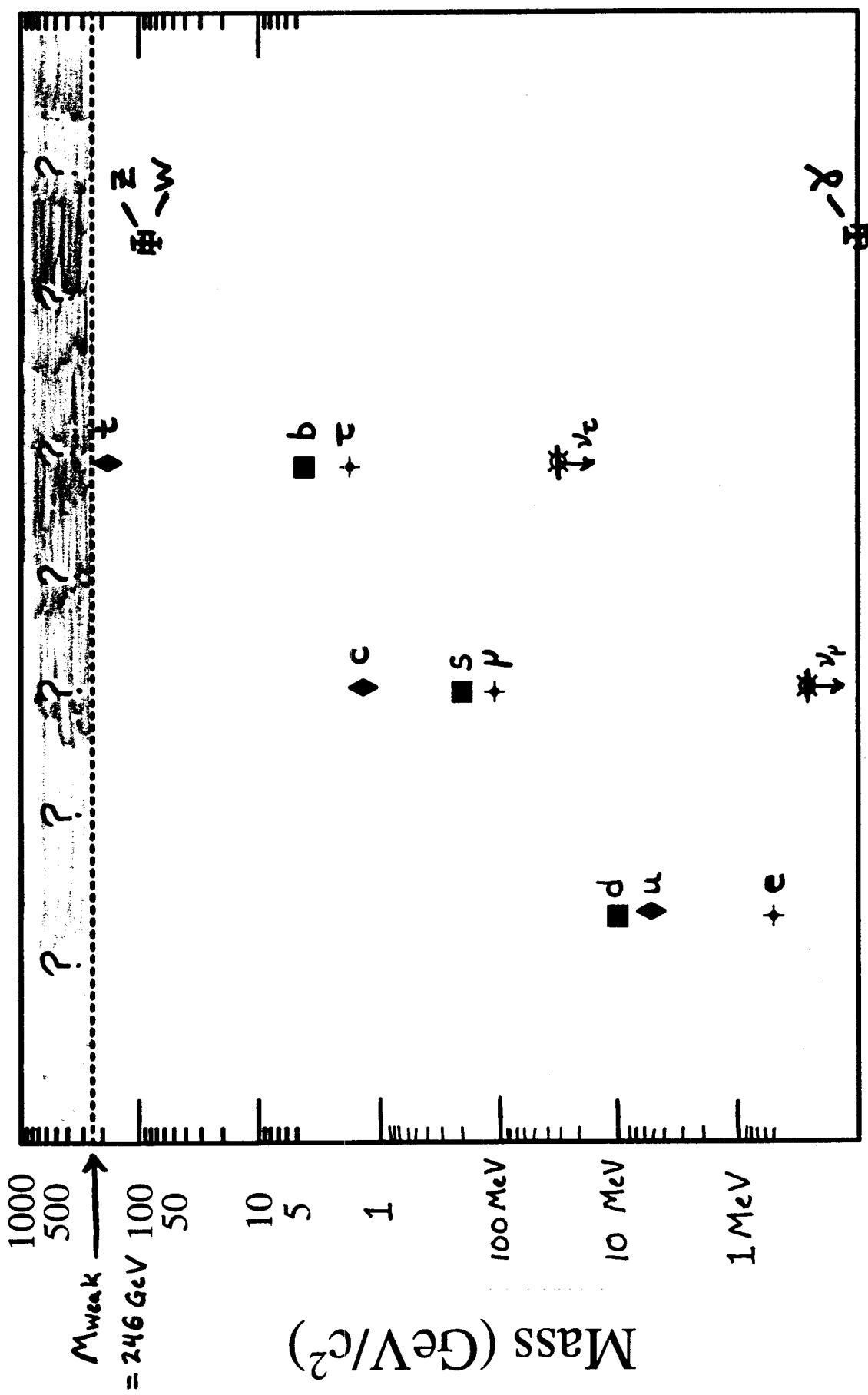
Despite continuing searches for superpartners we have not observed a "doubling" of the spectrum...



Should we be surprised?

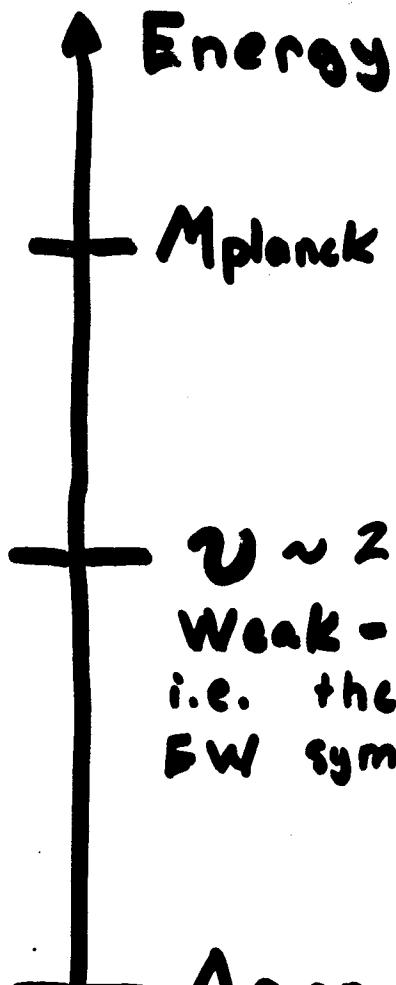
Can this continue?

Fundamental Particles



Fundamental Energy Scales

$$\hbar = c = 1$$



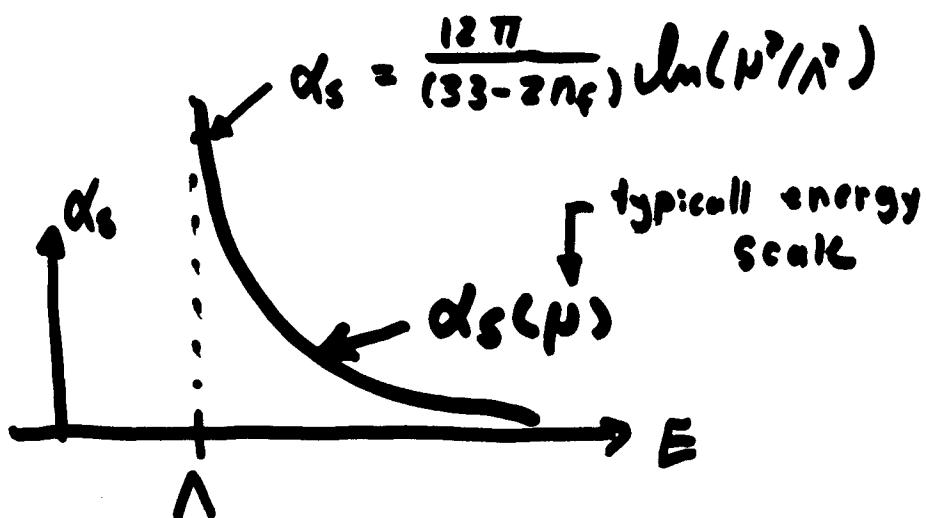
$$F_G \sim G_N \frac{m_1 m_2}{r^2}$$

$$G_N = \frac{1}{M_P^2}$$

$$G_F \sim \frac{1}{v^2} \approx 246 \text{ GeV}$$

$$M_W \sim g v$$

$$\langle \phi \rangle = v$$



All of the masses in the SM are related to the two fundamental scales v, Λ_{QCD} by dimensionless couplings

The Hierarchy Problem & The Naturalness Problem

M_{Pl}

$v = 246 \text{ GeV}$

Λ_{QCD}

Hierarchy Problem

why is $v \ll M_{Pl}$?

Naturalness Problem

If we impose $v \ll M_{Pl}$
is this stable.

By comparison $\Lambda_{\text{QCD}}/M_{Pl} \ll 1$ has no such problem:

$$\text{Hierarchy: } \frac{1}{g_s^2(\mu)} = \frac{1}{g_s^2(\mu')} + \frac{b}{(4\pi)^2} \ln\left(\frac{\mu'}{\mu}\right)$$

$$\Rightarrow \Lambda_{\text{QCD}} \sim M_{Pl} e^{-\frac{(4\pi)^2}{g_s^2(M_{Pl})}}$$

Naturalness: Symmetry = (broken) Scale invariance

If we change $g_s(M_{Pl})$ by a little bit

$\Lambda_{\text{QCD}}/M_{Pl}$ is still $\ll 1$

The Naturalness Problem in the SM

or why the SM can't be the correct theory
at energies $\gg v \sim 246 \text{ GeV}$

- Assume the SM is a good theory up to some scale Λ
- Compute the Higgs mass ($\alpha \text{ GeV}$)

$$m_H^2 = m_{0H}^2 + \delta m_H^2$$

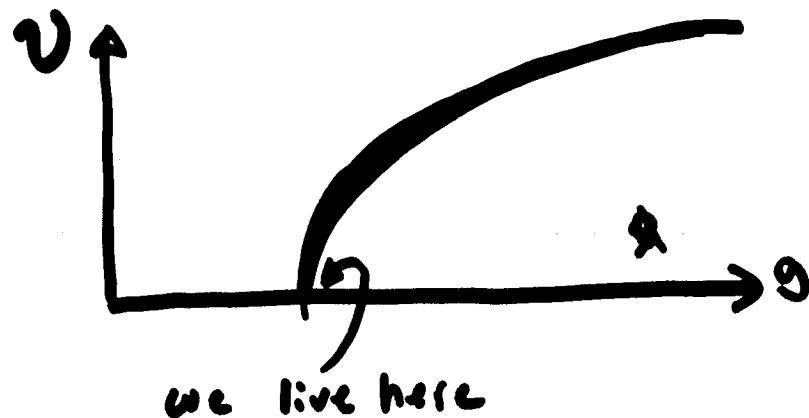
$\xrightarrow{\text{sum over int. states}}$
 $\approx \dots * \dots + \dots \underbrace{\dots}_{\text{at } E < \Lambda} \dots$

$$m_H^2 = m_{0H}^2 + g^2 \Lambda^2$$

Similarly

$$v^2 = -v_0^2 + g^2 \Lambda^2$$

Measured \downarrow \uparrow L e.g. $(10^{18} \text{ GeV})^2$
 $v^2 \approx (246 \text{ GeV})^2$ ~ 1

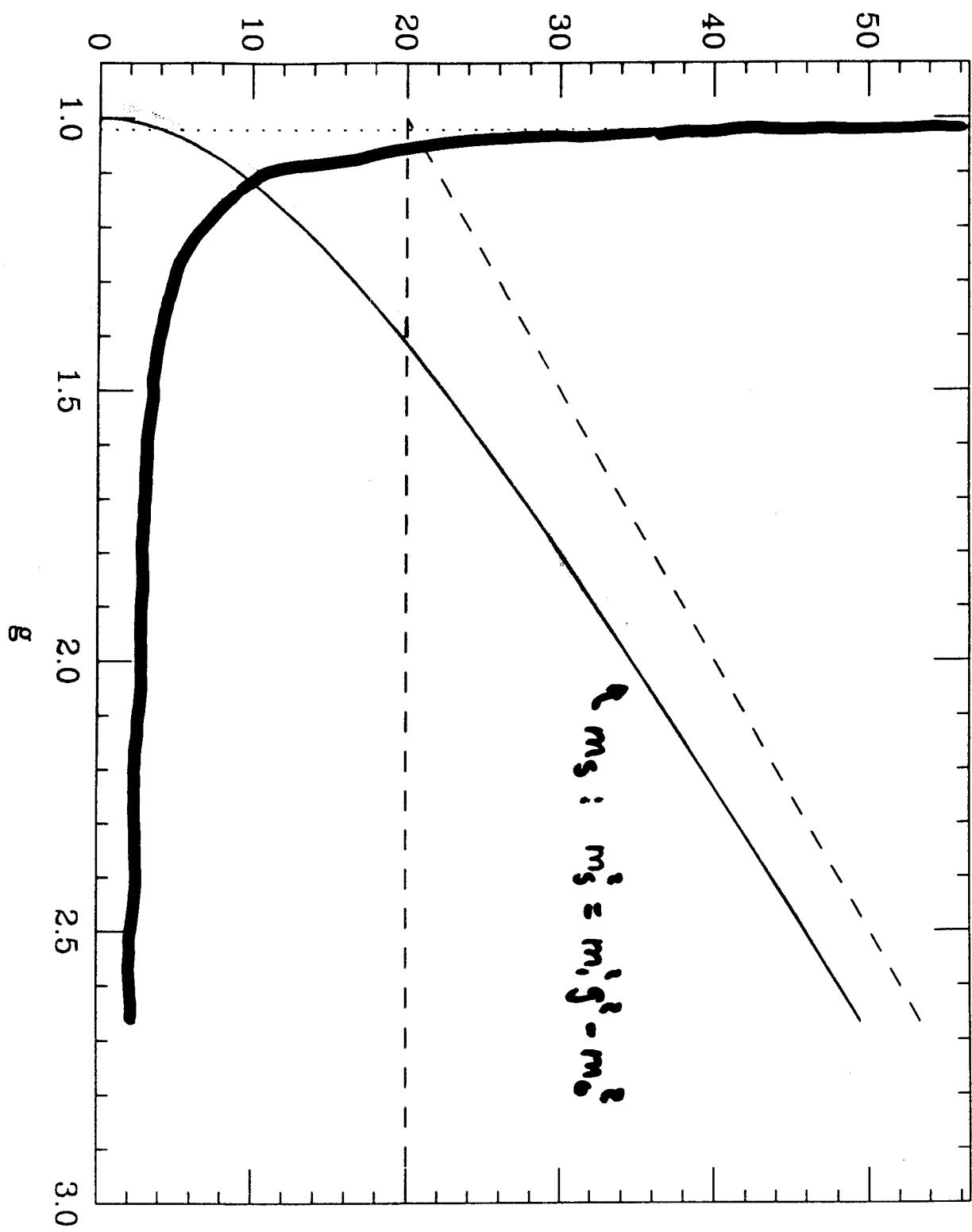


Requires

$$(246)^2 = -\underbrace{(10^{18})^2}_{\sim 1} + (10^{18})^2$$

Need to cancel
to 1 part in 10^{34}

Naturalness for a simple scalar



THE PHYSICS BEYOND THE SM
 i.e. at $E \gtrsim 0(100 \text{ GeV})$ can be
 classified by their solution to
 this problem. (why $\frac{v}{M_{\text{Pl}}} \ll 1$)

I Strong Dynamics ($S_{\text{ym}} = \text{Scale Inv}$)

The Higgs is not fundamental $\phi \sim f\bar{f}$
 v is from broken scale invariance
 like Λ_{QCD} in QCD

Examples : Technicolor, Topcolor

II Extra Dimensions

M_{Pl} is not what you think it is
 n - compact dim. of radius R

$$M_{\text{Pl}} = M_F (M_F R)^{n/2} \quad \text{if } M_F \sim 1 \text{ TeV}$$

implies $R = M_F (M_{\text{Pl}} / M_F)^{2/n}$



III Supersymmetry (Quarks...)

Scalars are fundamental

field theory works @ $E > 100 \text{ GeV}$

$$m_H^2 = -\dots \times \dots + \dots - \text{SM} \quad \downarrow \quad \text{SUPER PARTNER}$$

$+ g^2 \Lambda^2$ $- g^2 \Lambda^2$

3 Classic Naturalness Criteria

1937: DIRAC "No small numbers."

$$\lambda_i = \mathcal{O}(1)$$

$$\left. \begin{array}{l} \lambda_c = \mathcal{O}(10^{-5}) \\ \lambda_{\text{eff}}? \end{array} \right\}$$

1978: WILSON "Observable properties of a theory should be stable against minute variations of the fundamental parameters."

1980: 't HOOFT " $\alpha_i(\mu) \ll 1$ is only allowed if $\alpha_i(\mu) = c$ increases the symmetry of the system."



We will implement a modified form of the Wilson / Susskind criterion.

1995's Measurable quantities should not be unusually unstable against minute variations of the fundamental parameters.


S. Gherghetta

BEWARE OF THEORISTS BEARING LIMITS FROM NAT.

Unsuccessfull
attempts to quantify
naturalness

EENZ
BG

$$C = \frac{X}{Y} \cdot \frac{\delta Y}{\delta X}$$

↑ Major
Wilsonian
sensitivity
parameter

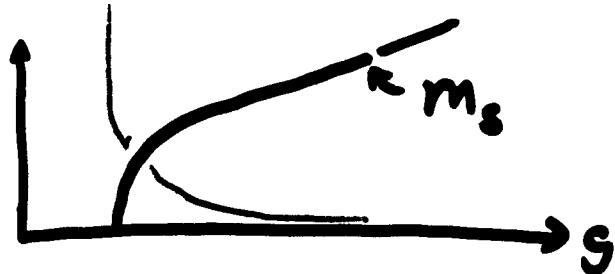
$$\frac{\delta y}{y} = C \frac{\delta x}{x}$$

y = observable

x = "fundamental" parameter

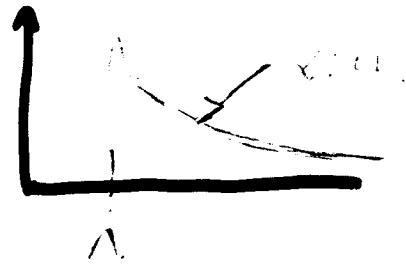
$$m_s^2 = m_\gamma g^2 - \Lambda^2$$

$$C = 2 \frac{m_\gamma g^2}{m_s^2}$$



So, why is this problematic?

$$M_{\text{prot}} \approx \Lambda_{\text{QCD}} = M e^{-\frac{8\pi^2}{b_2 g^2(\Lambda)}}$$



$$C = \frac{4\pi}{b_2} \left(\frac{1}{\alpha} \right) \gtrsim 100! \quad \frac{\delta M_p}{M_p} \approx (100) \frac{\delta \alpha}{\alpha}$$

EWSB $\frac{1}{2} m_Z^2 = \frac{m_{H_D}^2 - t_F^2 m_{H_U}^2}{t_F^2 - 1} - \mu^2$

Define a smooth metric on parameter space

$$dP = f(x) \frac{dx}{x} \\ = f(x) g(y) \frac{dy}{y}$$

probability that y lies within a scale invariant interval dy/g for a given x

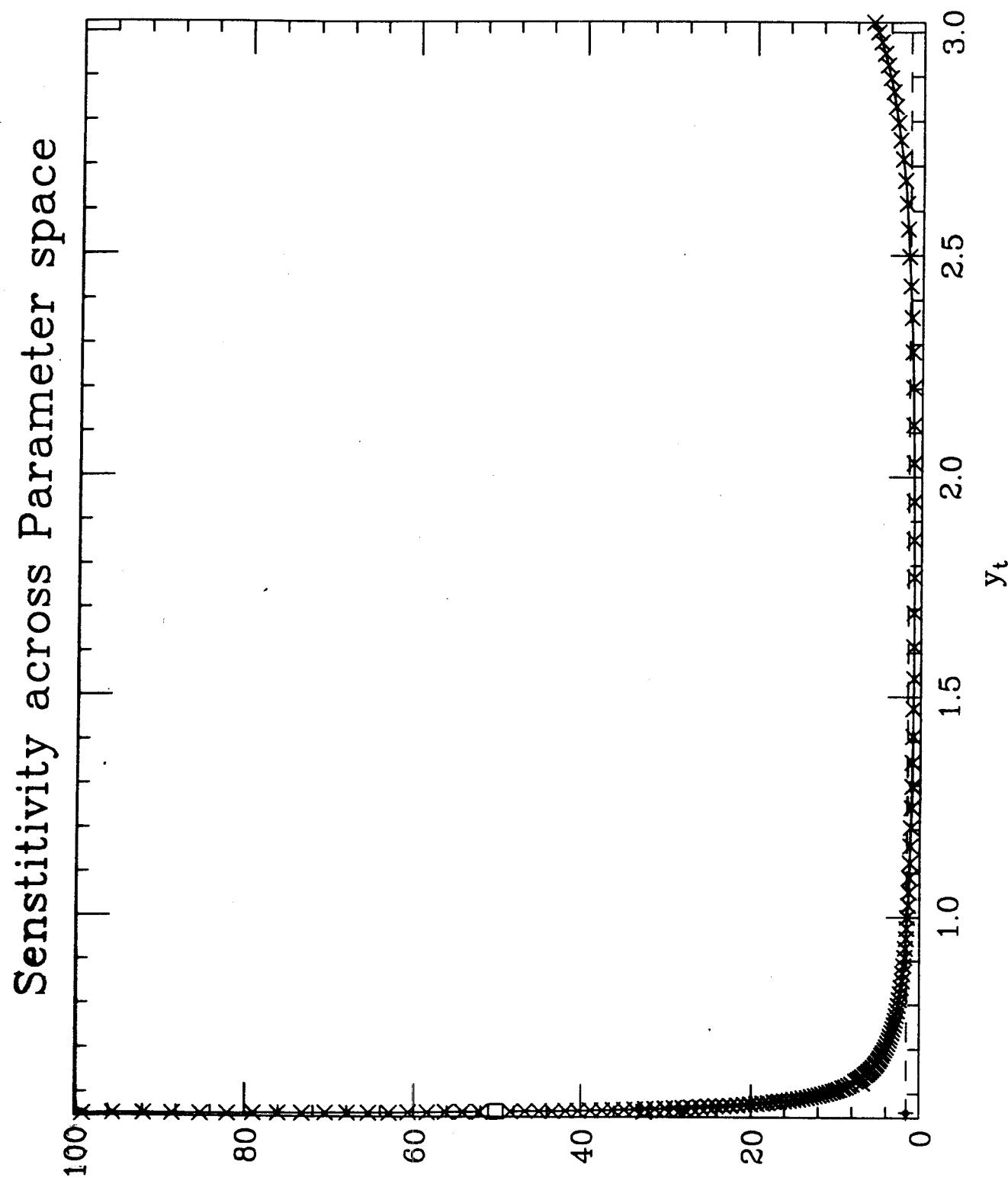
$$g(x) = \left[\frac{x}{y} \frac{dy}{dx} \right]^{-1}$$

Relative probability of observing y

$$\frac{g}{\bar{g}} = \frac{c^{-1}}{\int \frac{dx}{x} f c^{-1}}$$

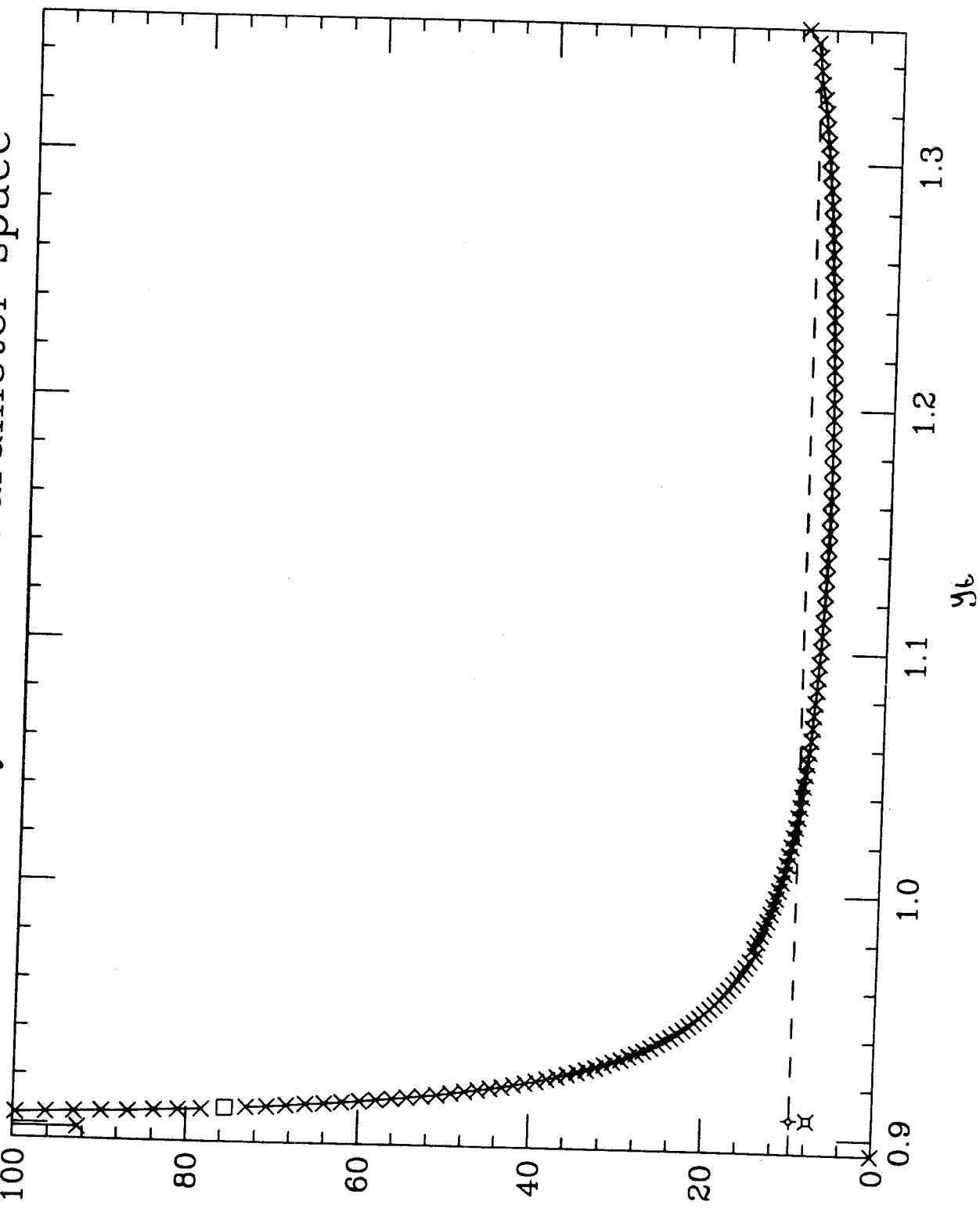
$$Y = \frac{\bar{g}}{g} = \frac{c}{\bar{c}}, \bar{c}^{-1} = -\frac{\int \frac{dx}{x} f c^{-1}}{\int \frac{dx}{x} f}$$

The observed value of y is fine-tuned when minute variations in fundamental parameters "x" lead to unusually large variations in y

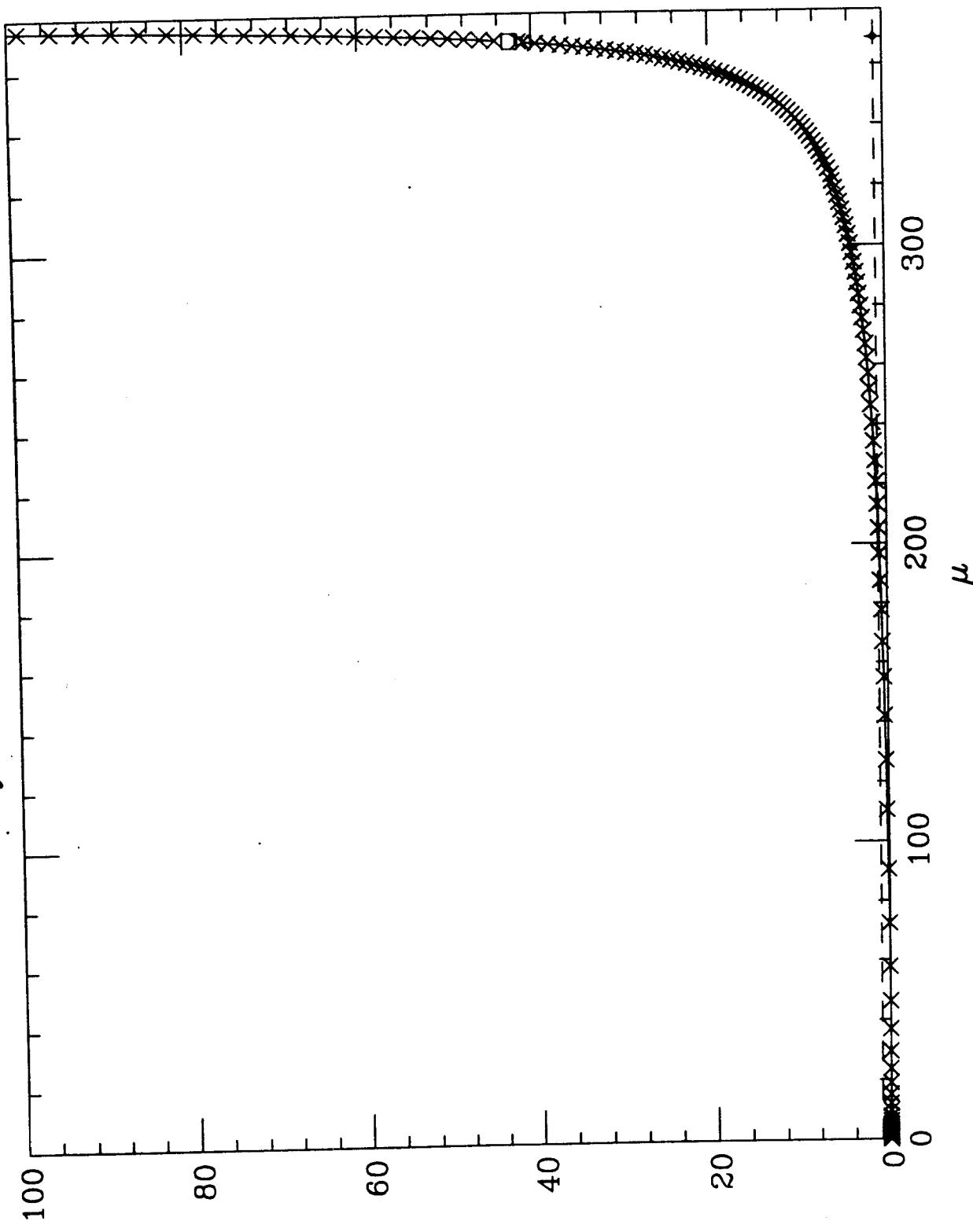


c

Sensitivity across Parameter space

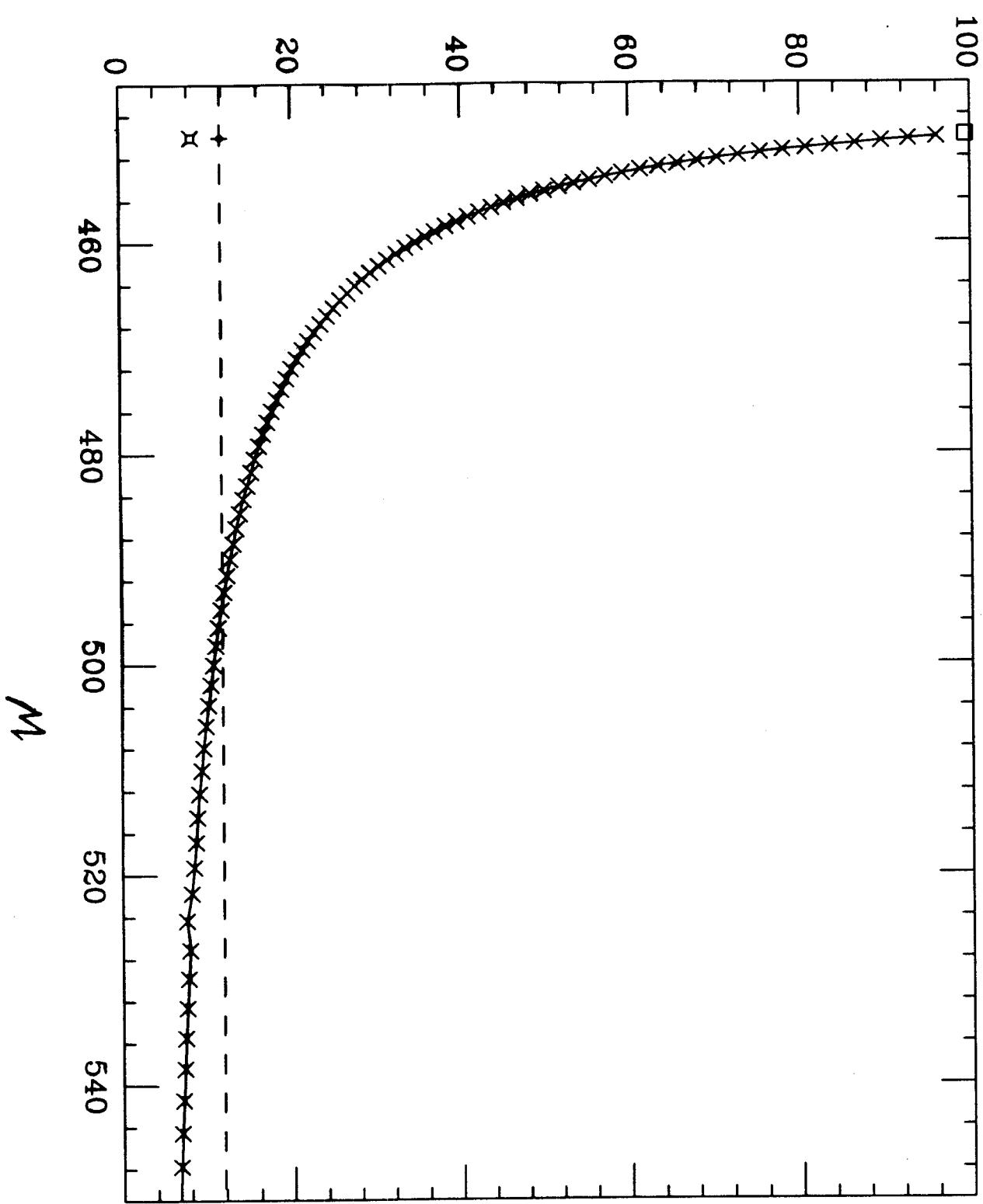


Sensitivity across Parameter space

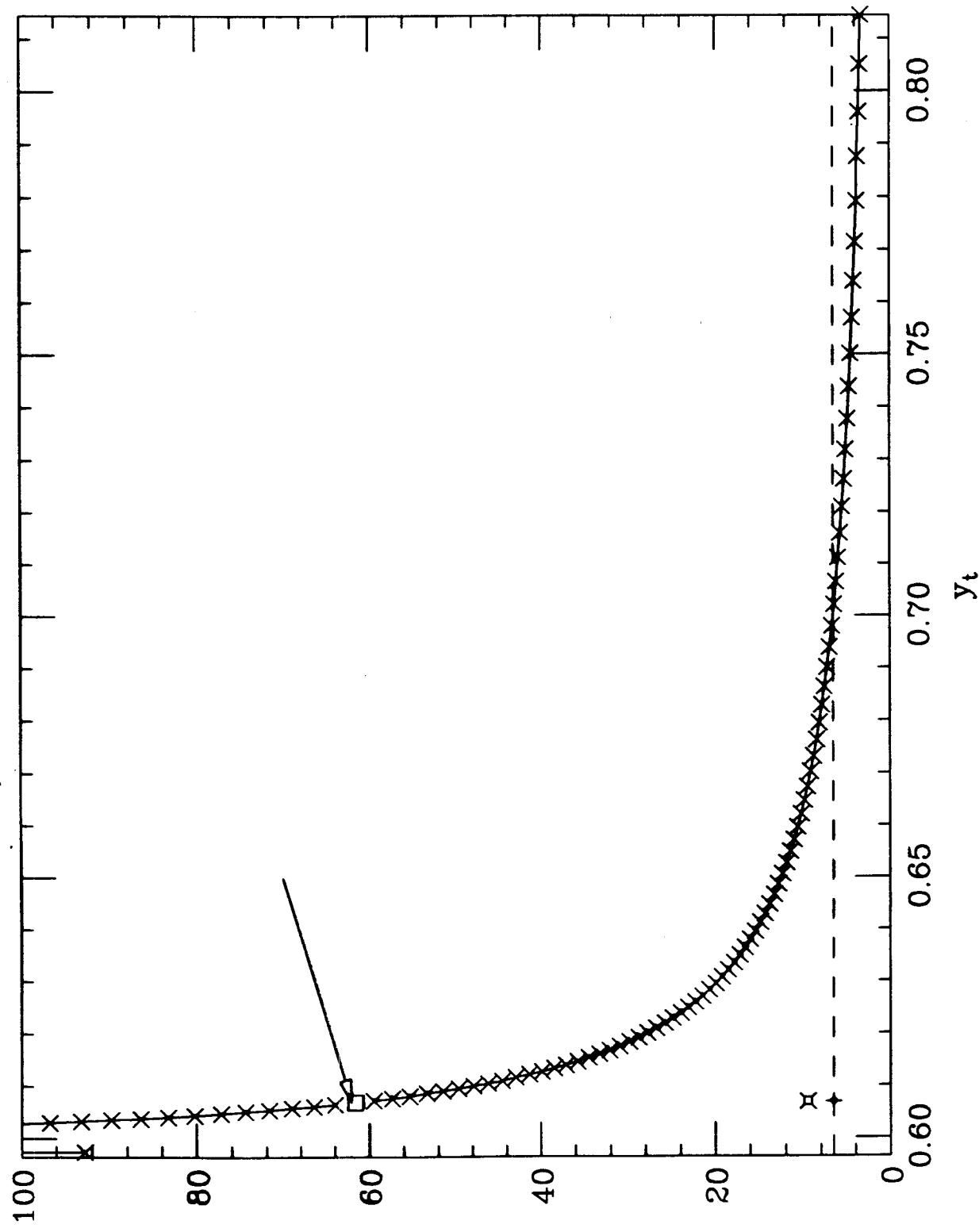


c

Sensitivity across Parameter space

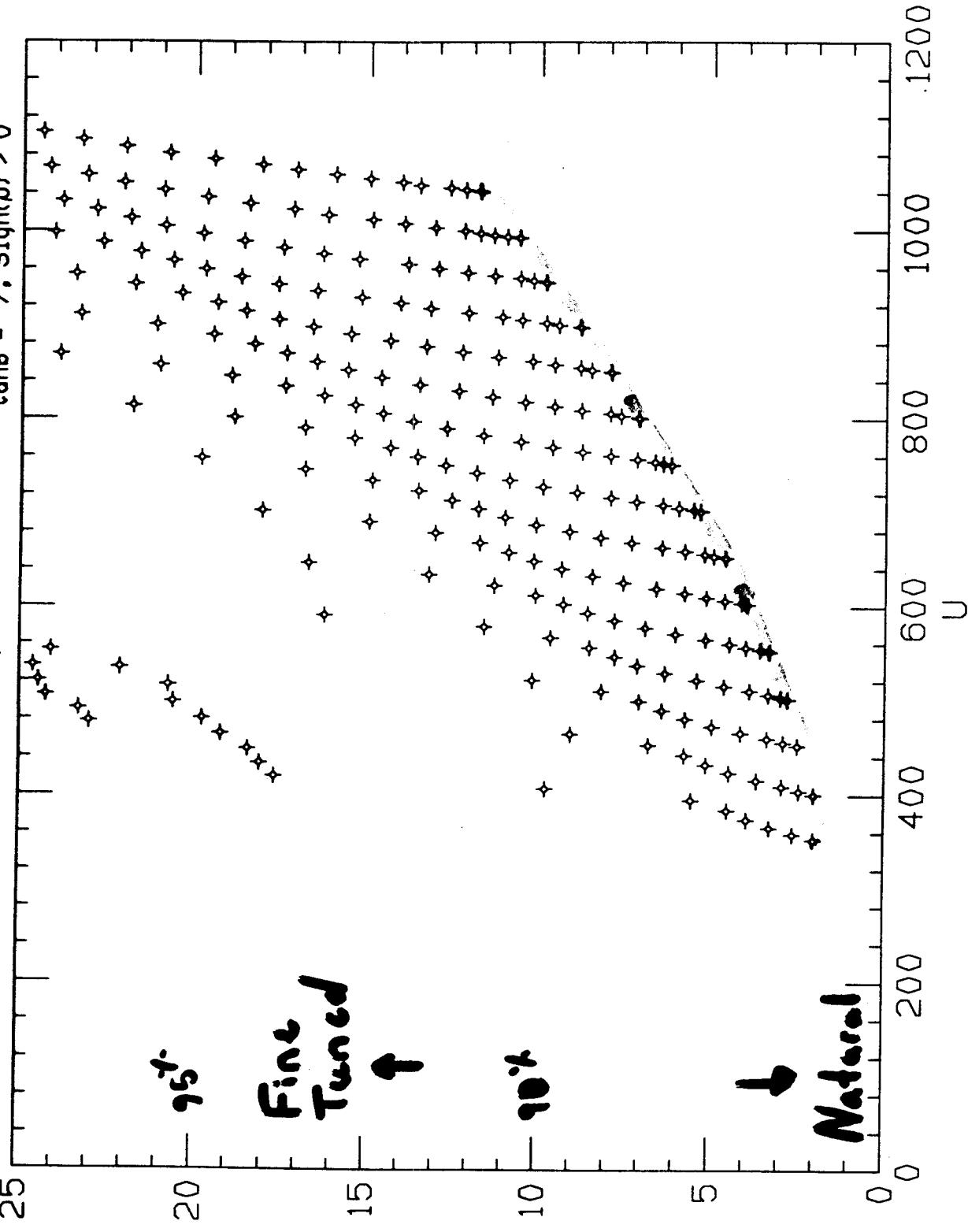


Sensitivity across Parameter space



c

Naturalness of Physical Particles vs. Their Masses



Y

Naturality of Physical Particles, Masses

25

20

15

10

8

5

0

st1

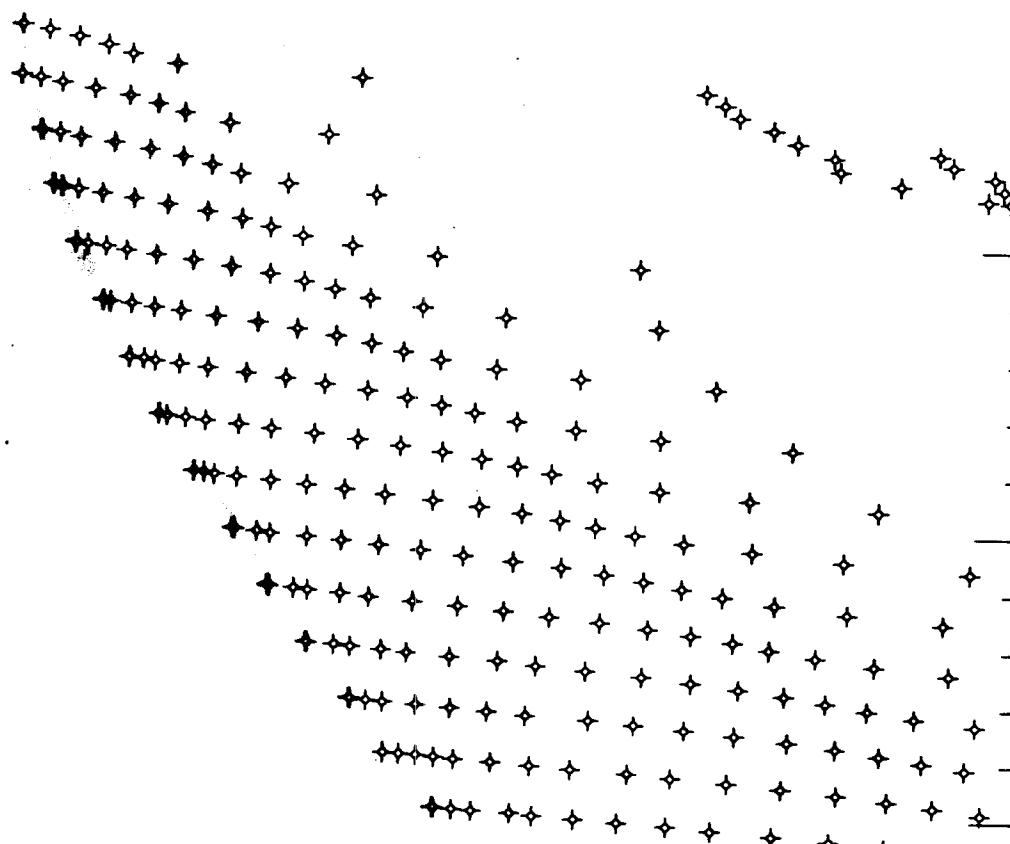
200

400

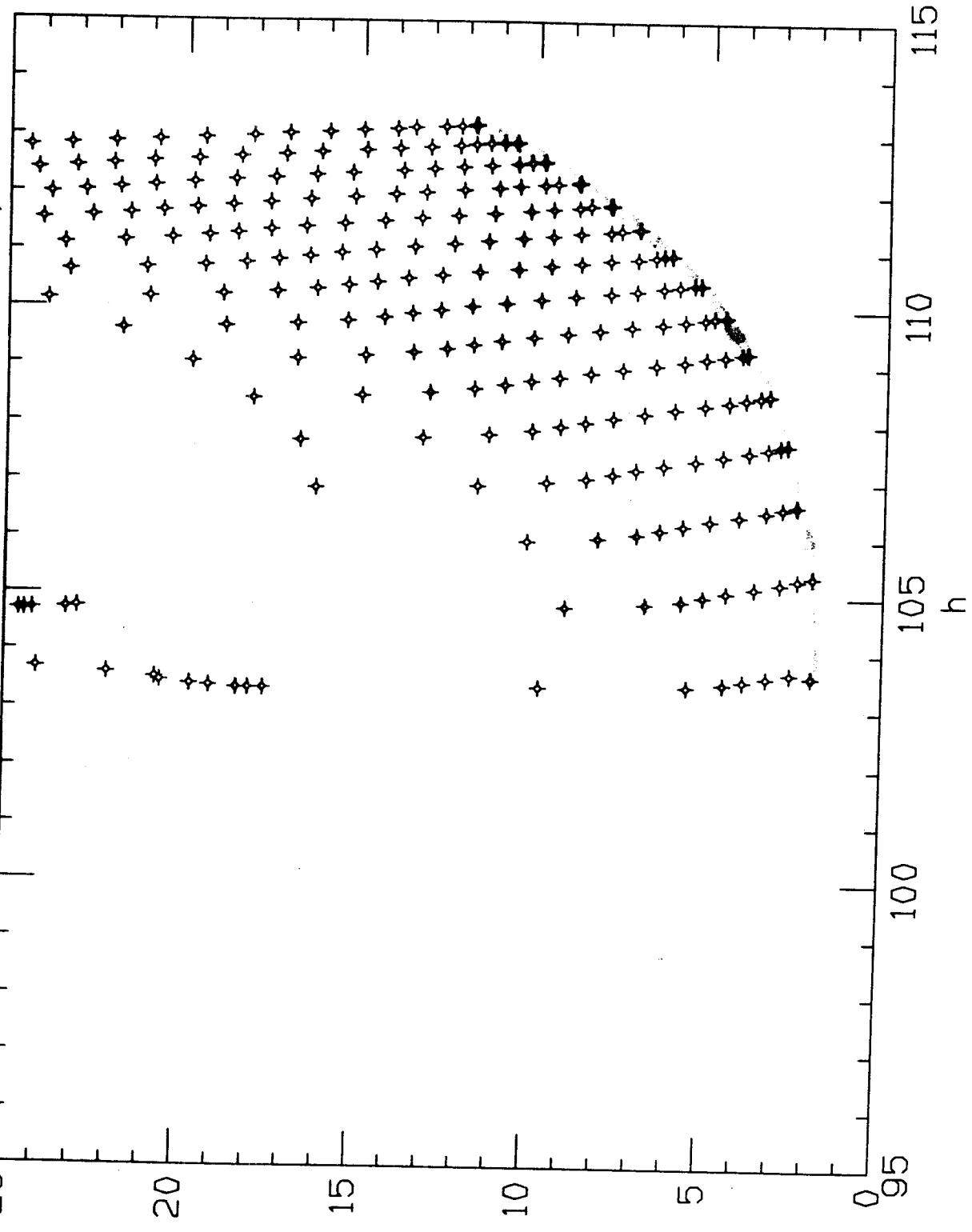
600

800

1000



Naturalness of Physical Particles, Signatures



Q

Naturality of Physical Particles Masses

25

20

15

8

10

5

0

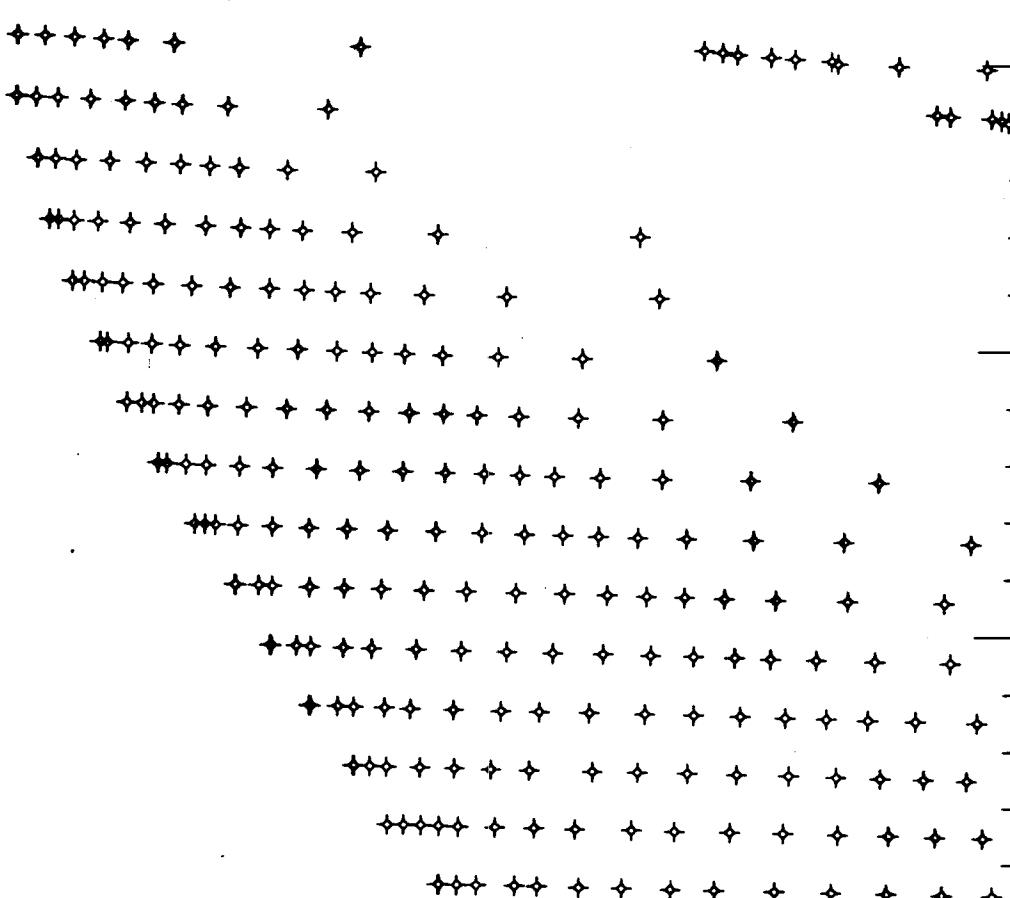
100

200

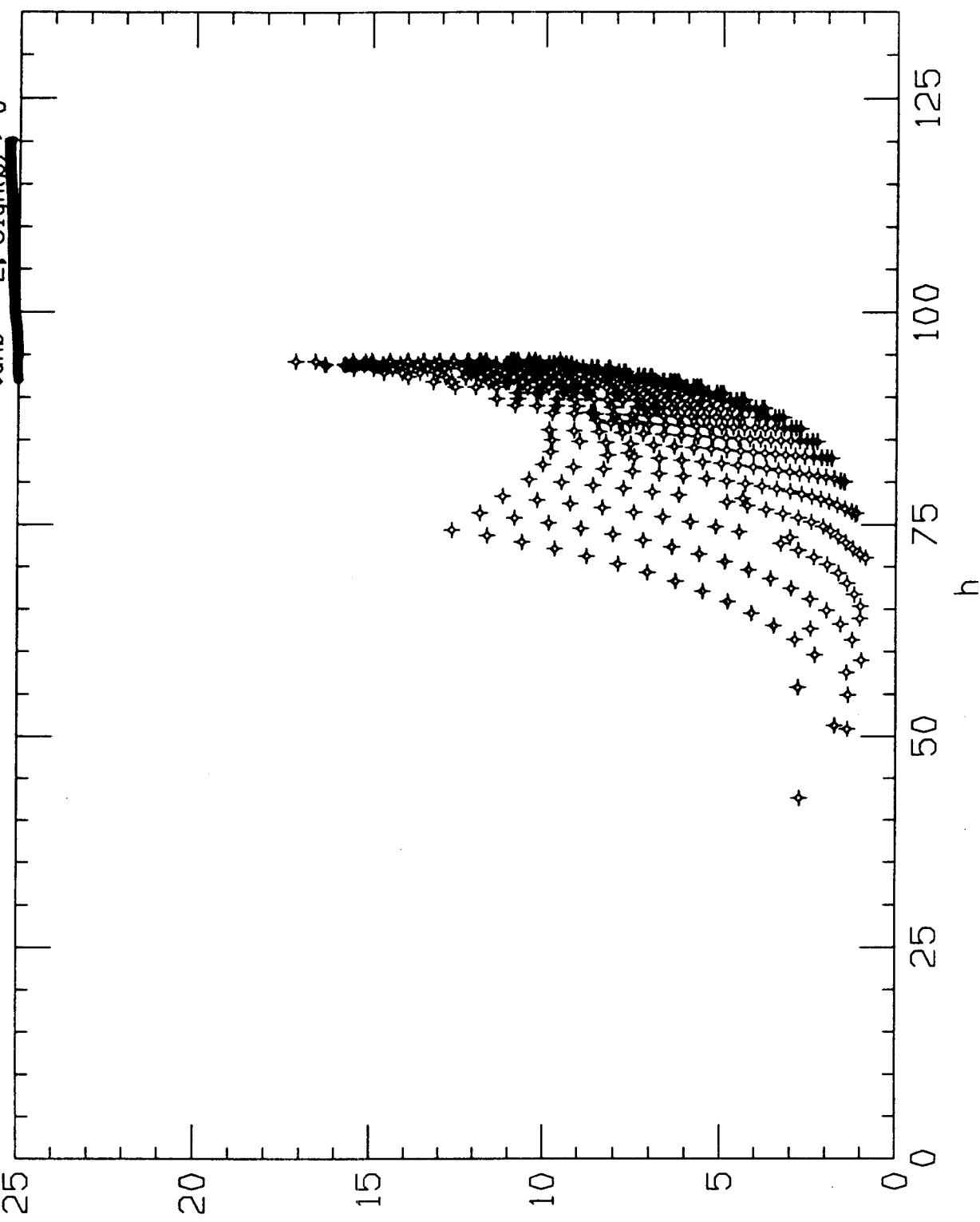
300

400

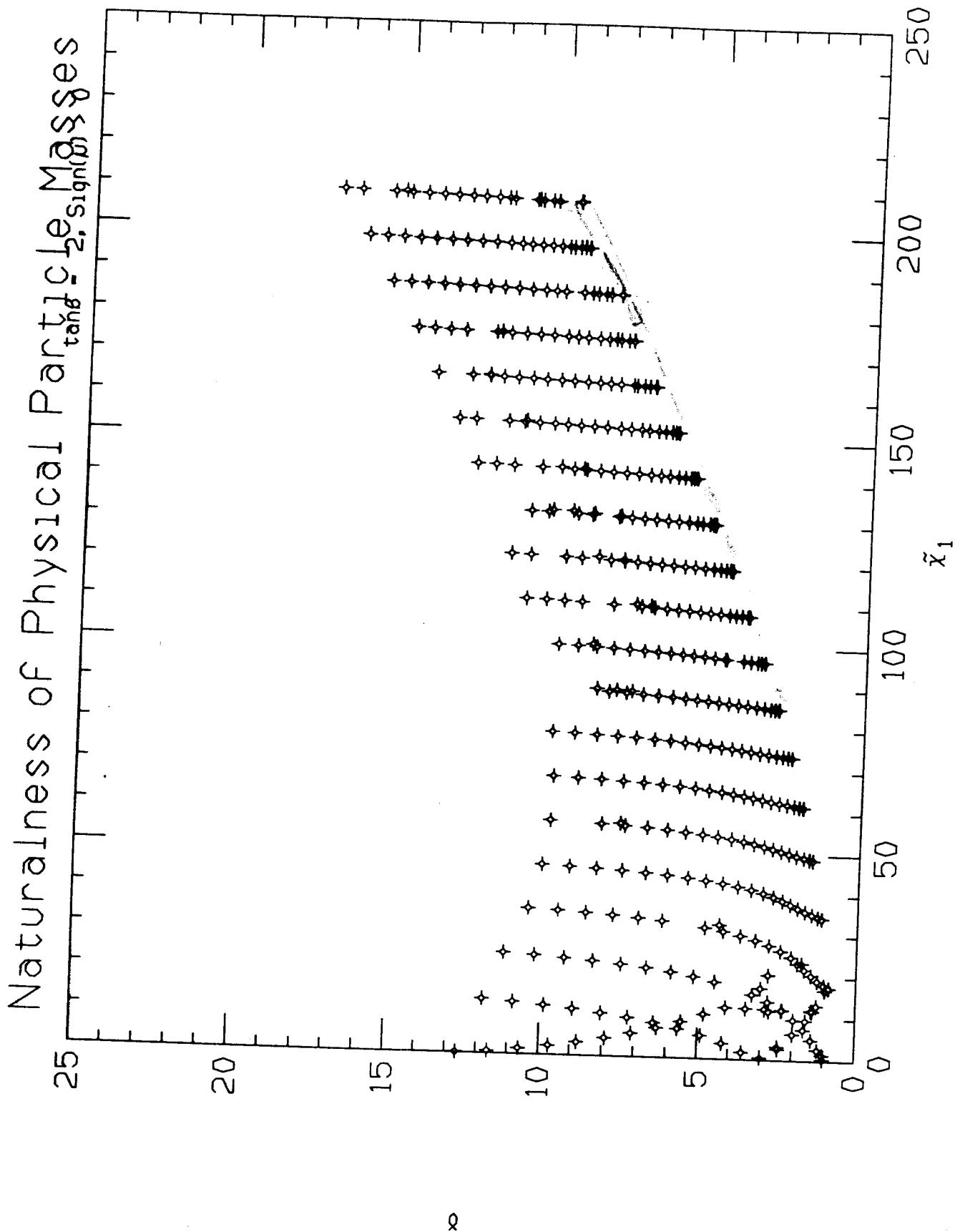
500



Naturalness of Physical Particles Masses

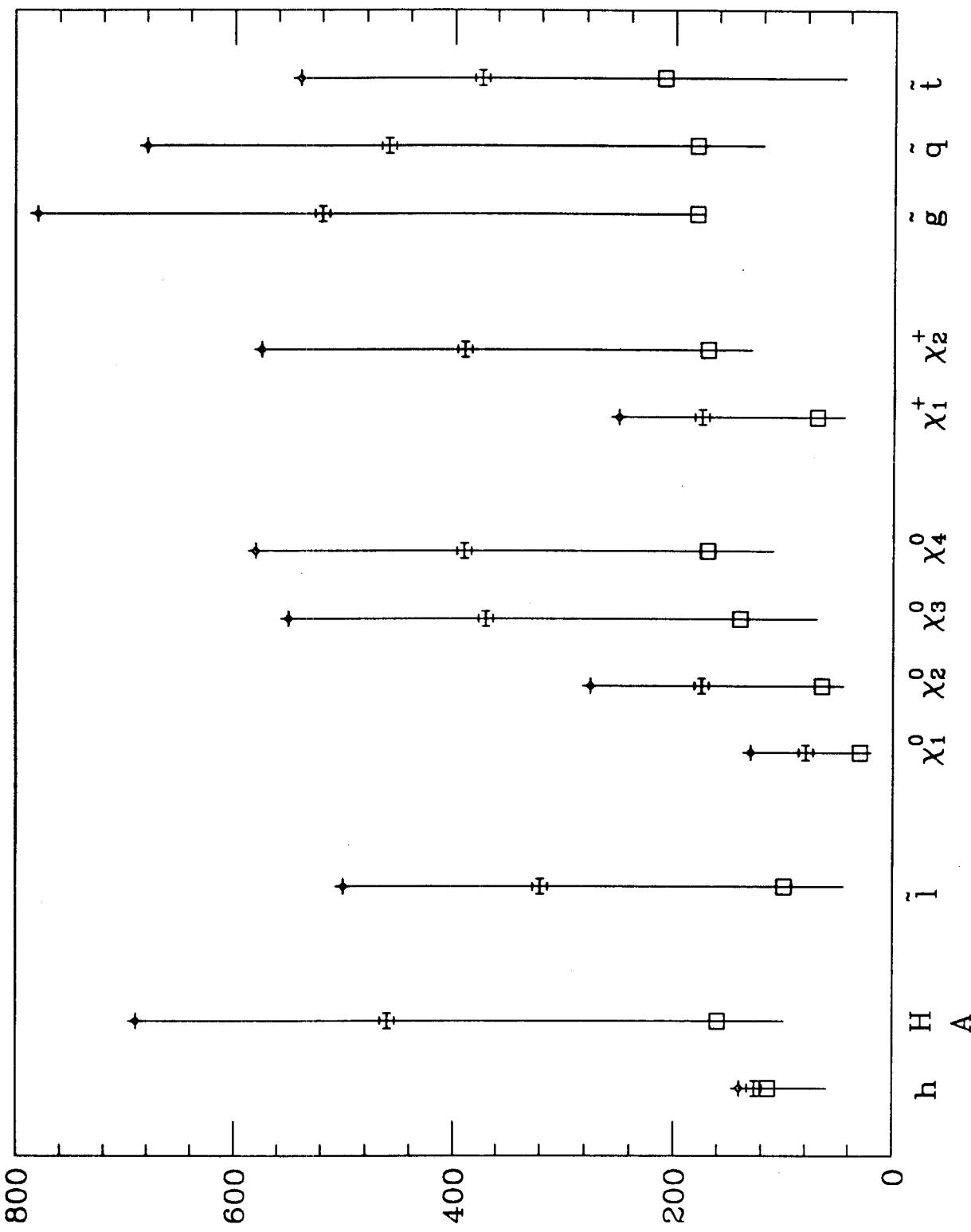


8



$\square \gamma = 1$
 $\times \gamma = 5$
 $\star \gamma = 10$

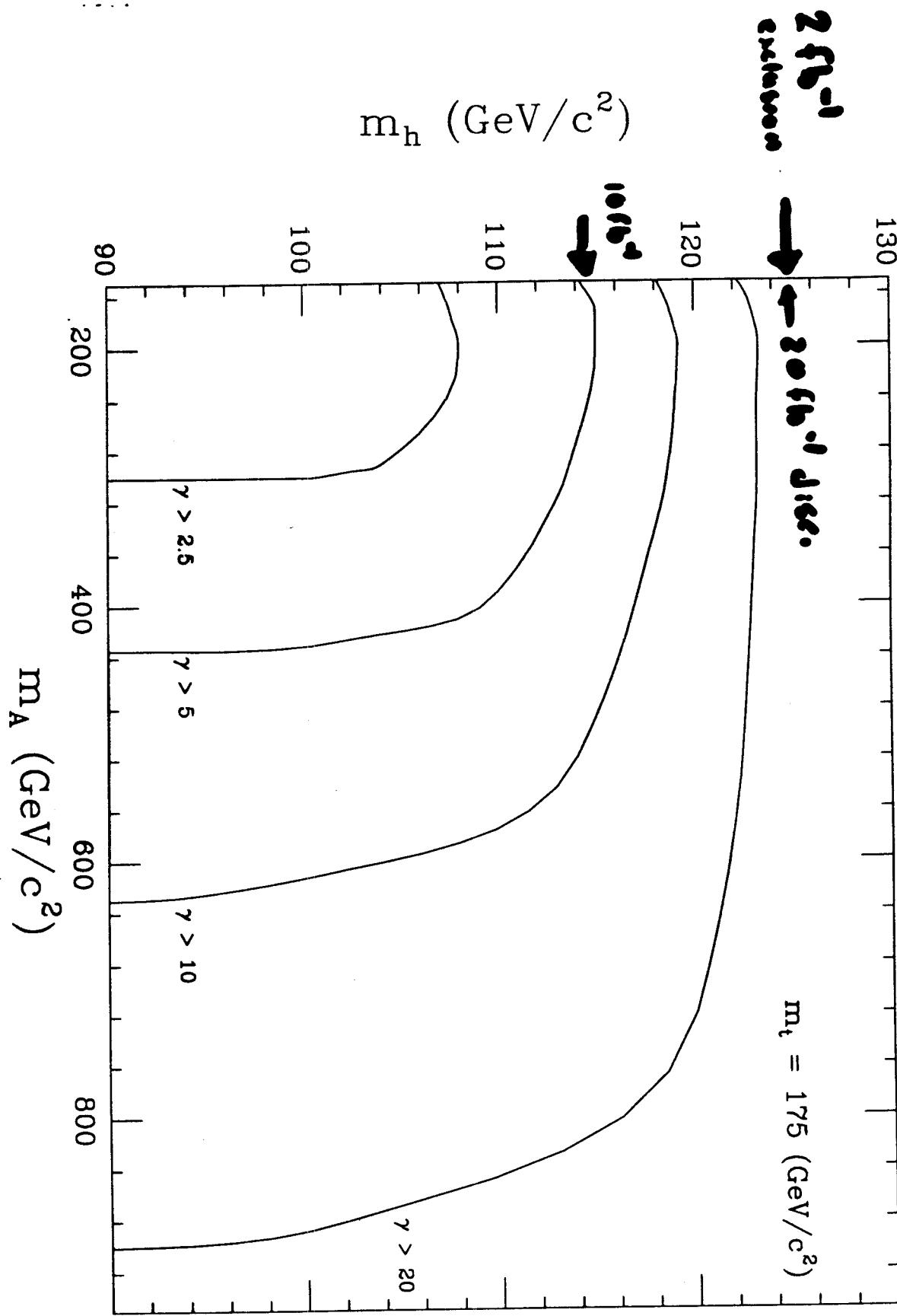
Naturalness and Superpartner Masses



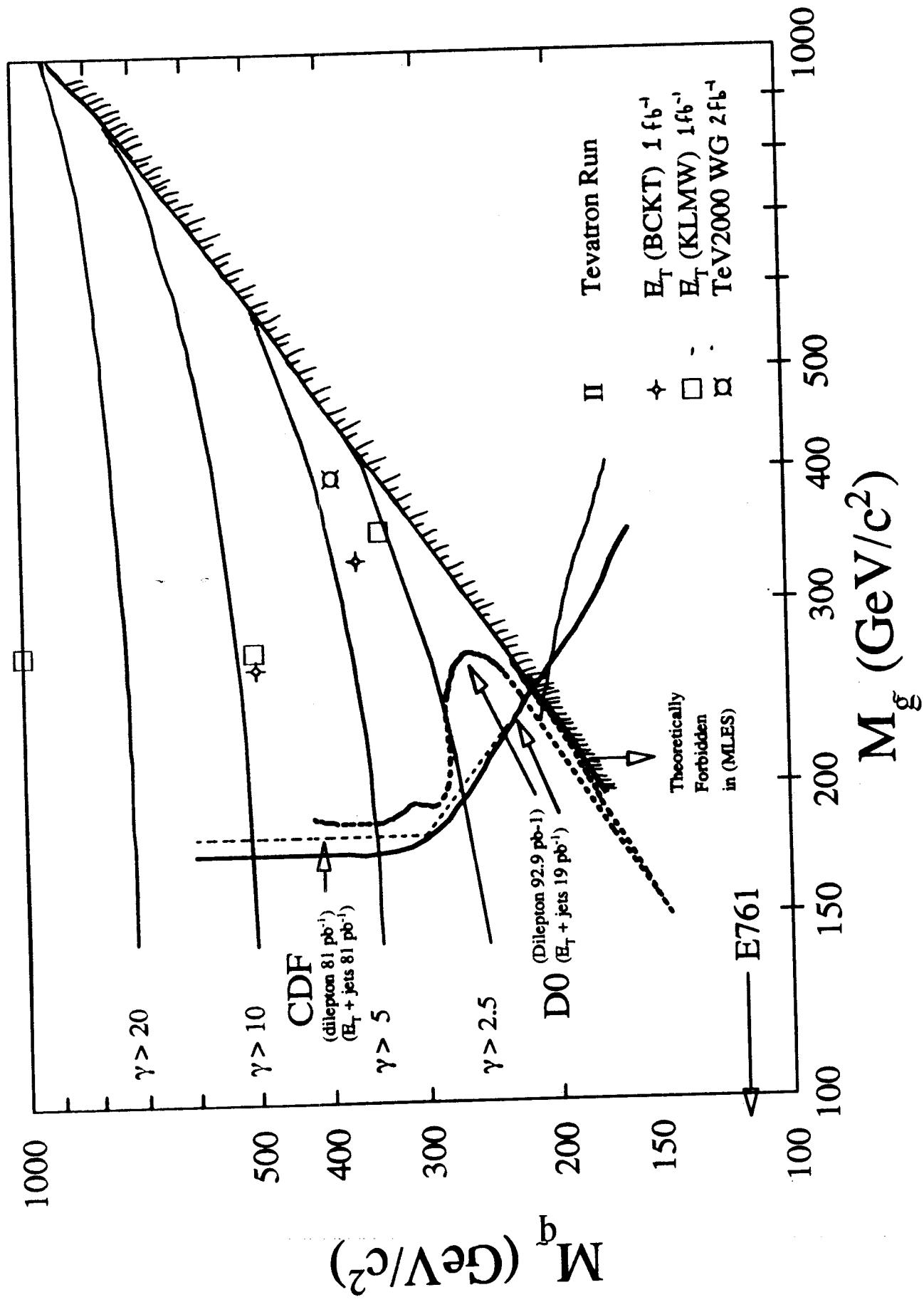
Mass (GeV/c²)

-
mg

Figure 2



Tevatron Searches



Conclusion and Outlook:

- If the next level beyond the SM is SUSY
 - All sparticles will be accessible to the LHC
 - The relevant scales are all $< TeV$
 - Beware of theorists with strongly time dependent naturalness limits

Tenable Scenarios for SUSY with new physics $> 1 \text{ TeV}$

- Scenarios motivating a scale $> M_{weak}$
 - Physics related to Extra Dimensions
 - Messengers
- Scenarios with add hoc New Scales
 - Heavy 1st & 2nd Generation Superpartners (FCNC)
 - Heavy Vectorlike reps